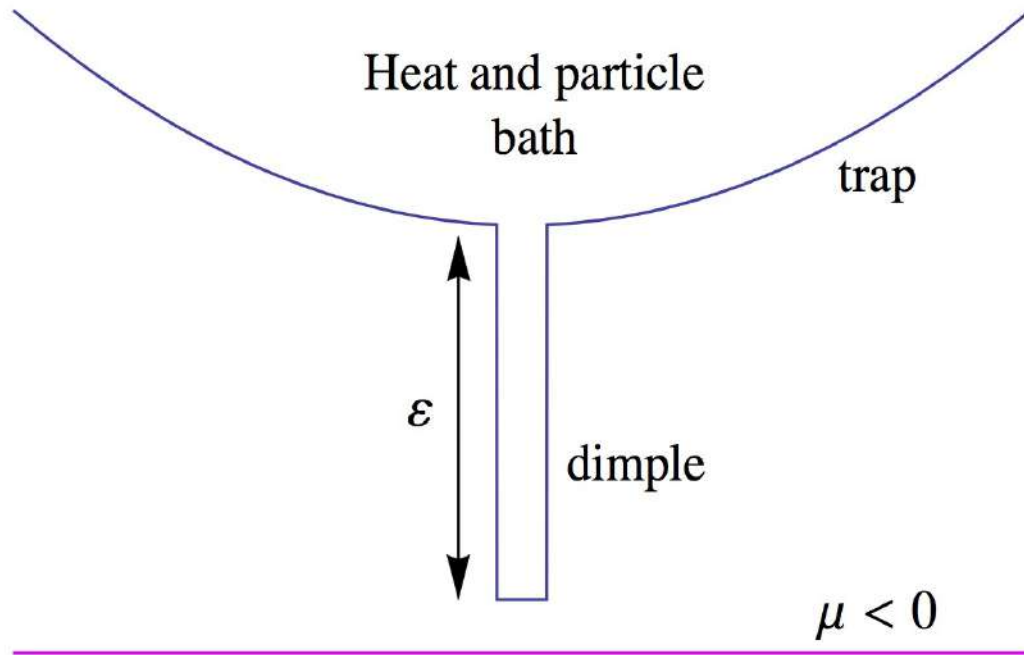


Kinetics of Bose condensation in a dimple potential

Shovan Dutta

Experimental system



Goal: Induce fast condensation in a narrow dimple potential by increasing particle density

Typical initial condition:

$$\mathcal{N} \approx 5 \times 10^8$$

$$T \approx 1 \mu\text{K}$$

$$\rho_{\text{bath}} \lambda_T^3 \approx 0.04 \ll 1$$

\implies bath is classical

$$\langle n_E \rangle = \frac{1}{e^{\beta(E-\mu)} - 1}$$
$$\approx e^{-\beta(E-\mu)}$$

Model

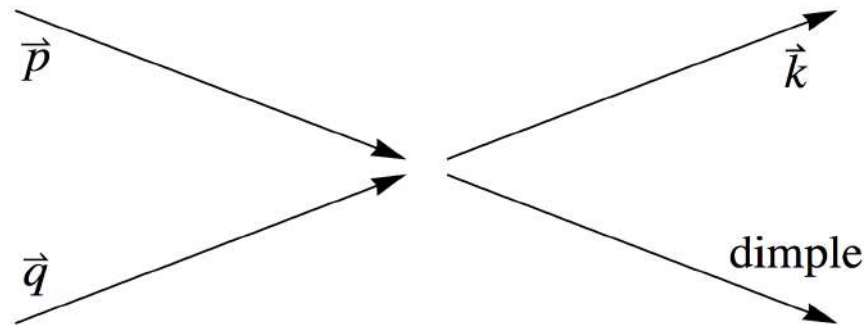
- ★ The bath: classical Boltzmann gas in a harmonic well having a thermal distribution
- ★ The dimple: quantum Bose gas in a square well
- ★ The bath sets T and μ
- ★ Rough idea: assuming thermal equilibrium between bath and dimple,

$$\langle n_E \rangle = \frac{1}{e^{\beta(E-\mu)} - 1}$$

\implies for condensation, $\mu \approx -\varepsilon_{\text{dimple}}$

Infinite trap: kinetics

★ Two-body collision



★ Rate proportional to

i) occupation: $e^{-\beta(\frac{p^2+q^2}{2m}-2\mu)}$

ii) Bose stimulation: $1 + N_{\vec{n}}$

iii) Fermi's golden rule: $\frac{2\pi}{\hbar} \left(\frac{4\pi\hbar^2 a_s}{m} \right)^2 |\tilde{\psi}_{\vec{n}}(\vec{p} + \vec{q} - \vec{k})|^2$

Infinite trap: kinetics

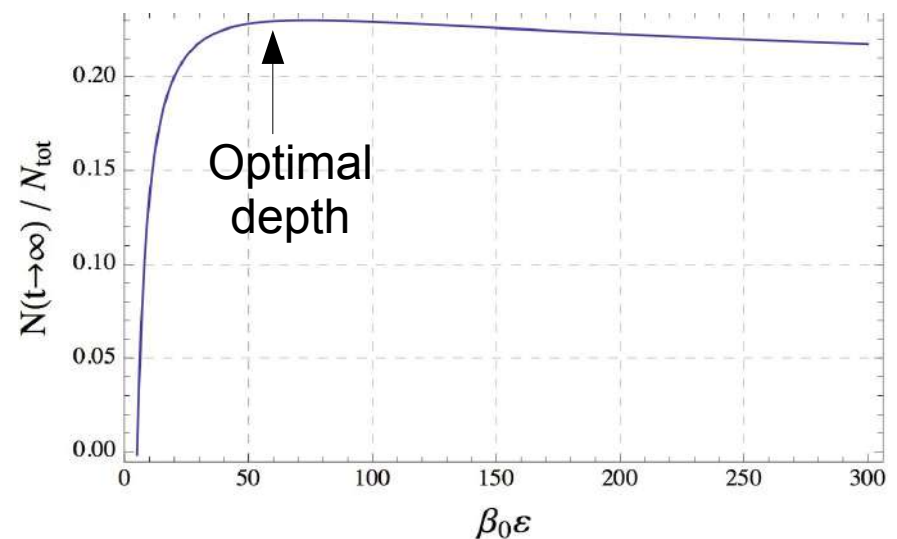
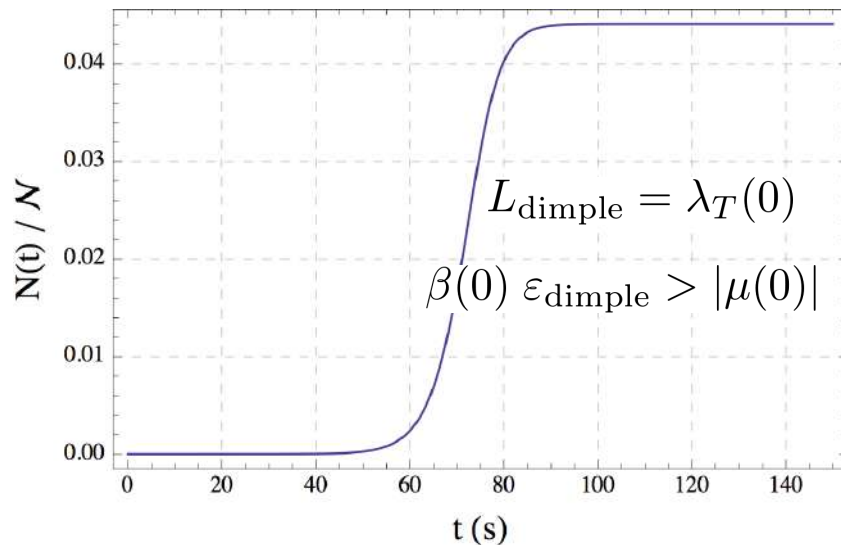
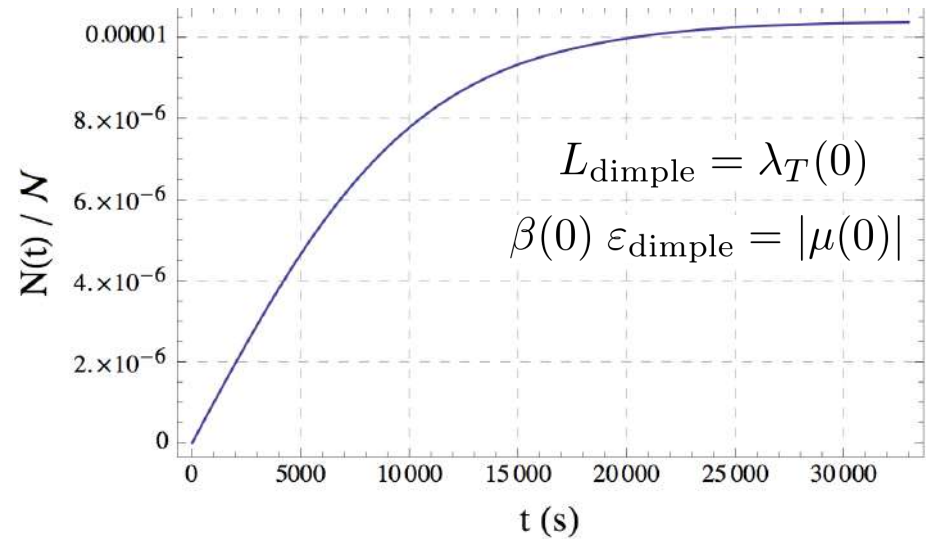
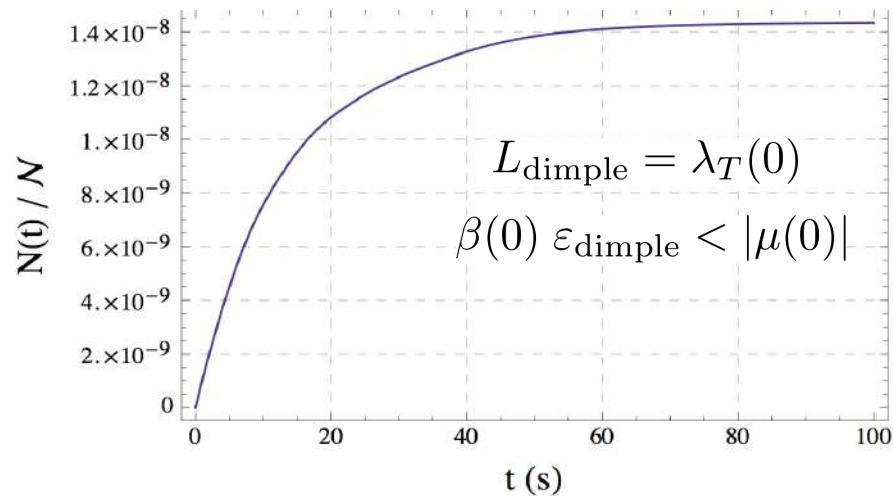
$$\psi_{\vec{n}}(\vec{r}) = \frac{1}{L_{\text{dimple}}^{3/2}} e^{i \frac{2\pi}{L_{\text{dimple}}} \vec{n} \cdot \vec{r}} ; \quad \varepsilon_{\vec{n}} = -\varepsilon_{\text{dimple}} + \frac{2\pi^2 \hbar^2}{m L_{\text{dimple}}^2} \vec{n}^2$$

$$\left(\frac{dN_{\vec{n}}}{dt} \right)_{in} = \frac{2\pi}{\hbar} \left(\frac{4\pi \hbar^2 a_s}{m} \right)^2 \int \frac{d^3 p d^3 q}{(2\pi \hbar)^6} e^{-\beta \left(\frac{p^2 + q^2}{2m} - 2\mu \right)} (1 + N_{\vec{n}}) \\ \times \delta \left((p^2 + q^2 - (\vec{p} + \vec{q} - 2\pi \vec{n} / L_{\text{dimple}})^2) / 2m + \varepsilon_{\vec{n}} \right)$$

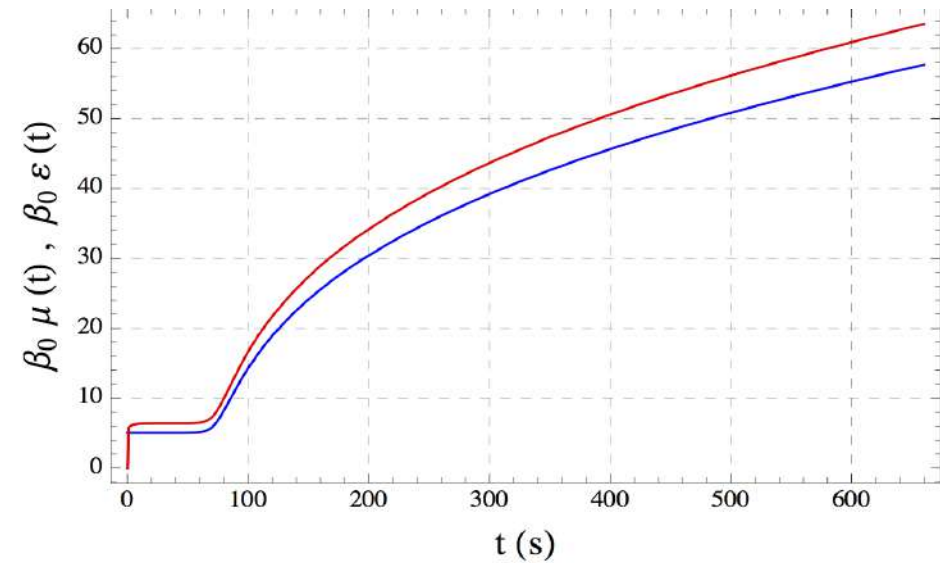
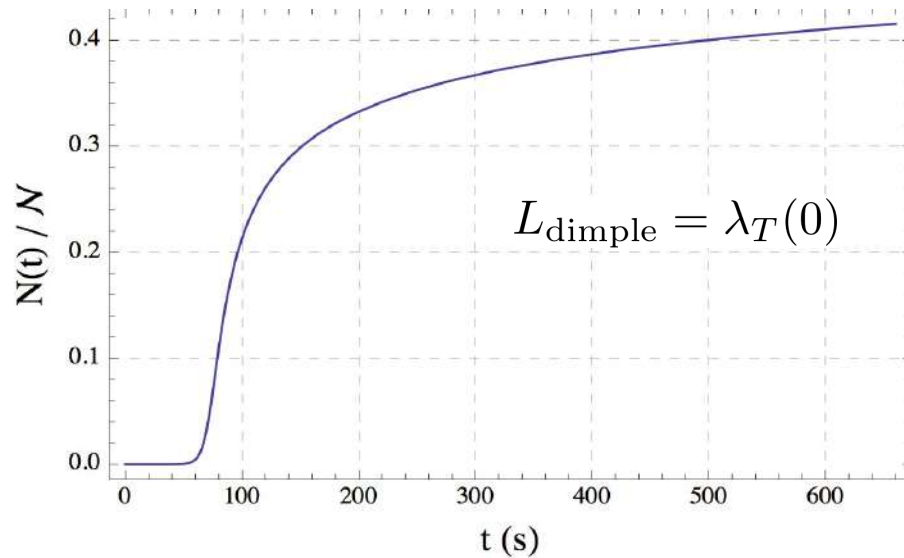
$$\left(\frac{dN_{\vec{n}}}{dt} \right)_{out} = -\frac{2\pi}{\hbar} \left(\frac{4\pi \hbar^2 a_s}{m} \right)^2 \int \frac{d^3 p d^3 q}{(2\pi \hbar)^6} e^{-\beta \left(\frac{k^2}{2m} - \mu \right)} N_{\vec{n}} \\ \times \delta \left((p^2 + q^2 - (\vec{p} + \vec{q} - 2\pi \vec{n} / L_{\text{dimple}})^2) / 2m + \varepsilon_{\vec{n}} \right)$$

$$N_{\text{bath}}(t) = \mathcal{N} - N_{\text{dimple}}(t) ; \quad \dot{E}_{\text{bath}}(t) = -\varepsilon_{\vec{n}}(t) \dot{N}_{\text{dimple}}(t)$$

Time evolution for one bound state



Optimal evolution for one bound state



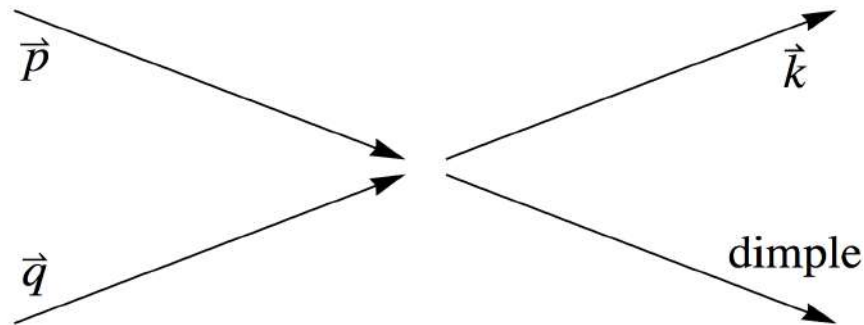
★ However, in typical experiments, $L_{\text{dimple}} \gg \lambda_T(0)$ and

$$\beta(0) \epsilon_{\text{dimple}} \gtrsim 5$$

$$\implies \text{number of bound states} \approx \sqrt{\frac{\beta \epsilon_{\text{dimple}}}{\pi}} \frac{L_{\text{dimple}}}{\lambda_T} \gg 1$$

Many bound states: interaction kinetics

★ Transition between two energy levels:



★ Rate $\propto \exp[-\beta(p^2/2m - \mu)] N_{\vec{n}}$ and $1 + N_{\vec{m}}$.

$$\begin{aligned} \frac{dN_{\vec{n} \rightarrow \vec{m}}}{dt} &= \frac{2\pi}{\hbar} \left(\frac{4\pi\hbar^2 a_s}{m} \right)^2 \left(\frac{2\pi\hbar}{L_{\text{dimple}}} \right)^3 \int \frac{d^3 p d^3 k}{(2\pi\hbar)^6} e^{-\beta(\frac{p^2}{2m} - \mu)} \\ &\times N_{\vec{n}} (1 + N_{\vec{m}}) \delta(p^2/2m + \varepsilon_{\vec{n}} - k^2/2m - \varepsilon_{\vec{m}}) \\ &\times \delta^3(\vec{p} + 2\pi\vec{n}/L_{\text{dimple}} - \vec{k} - 2\pi\vec{m}/L_{\text{dimple}}) \end{aligned}$$

Many bound states: continuum approximation

$$\sum_{\vec{n}} N_{\vec{n}}(t) \longrightarrow \mathcal{N} \int f(\beta(0)E, t) d(\beta(0)E)$$

$$f(\beta(0)E, t) = g(\beta(0)E_{\vec{n}}) N_{\vec{n}}(t) / \mathcal{N}$$

where $g(\beta(0)E) = 2 (L_{\text{dimple}} / \lambda_T(0))^3 \sqrt{\beta(0)E / \pi}$ is the density of states from the bottom of the dimple.

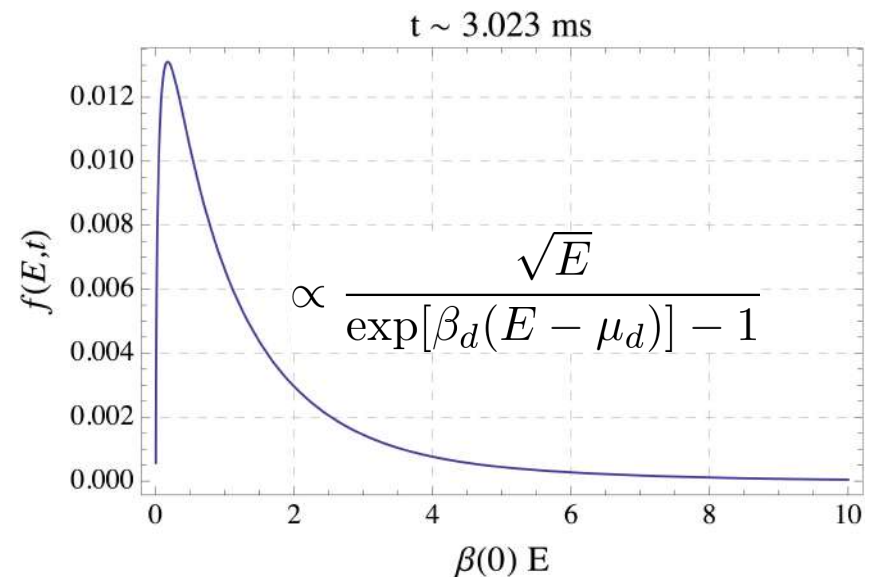
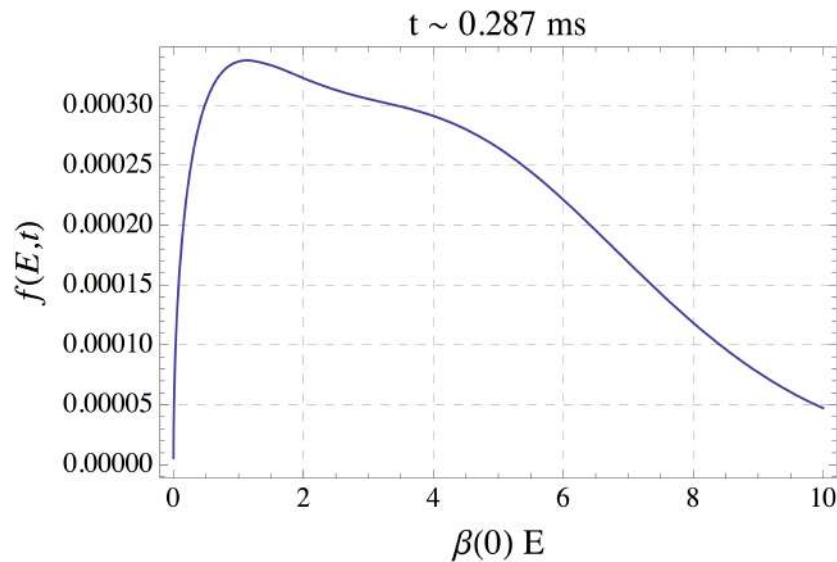
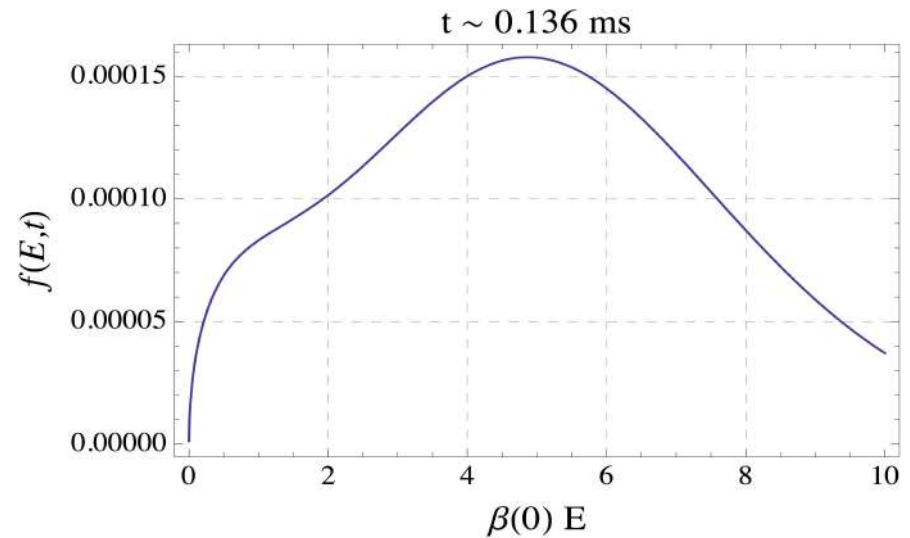
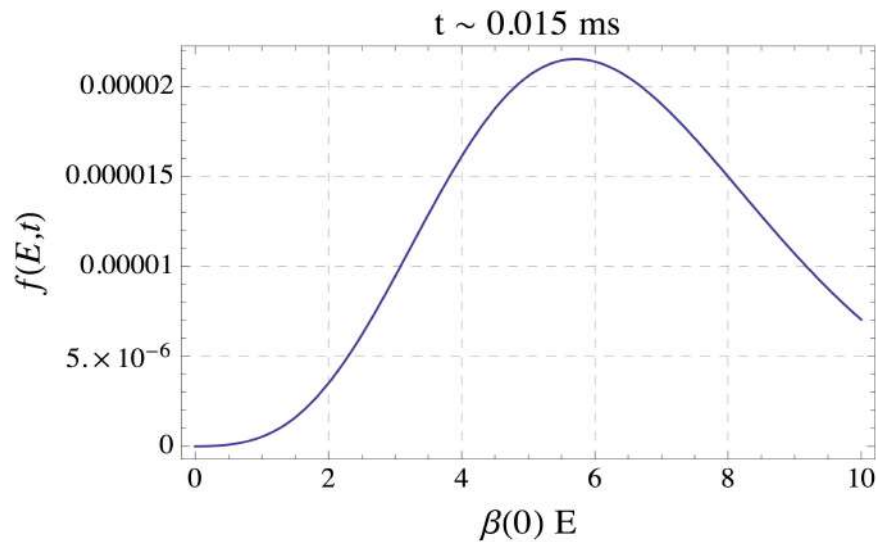
- ★ Consider the zero energy state separately to observe condensation.
- ★ Condensation occurs when $\rho_{\text{bath}}(0) \lambda_T^3(0) (V_{\text{bath}}(0) / V_{\text{dimple}}) \gg 1$ as expected.

Interaction causes thermal equilibrium within dimple

$$\beta(0) \varepsilon_{\text{dimple}} = 10$$

$$\rho_{\text{bath}}(0) \lambda_T^3(0) \approx 0.05$$

$$V_{\text{bath}}(0)/V_{\text{dimple}} = 2000$$

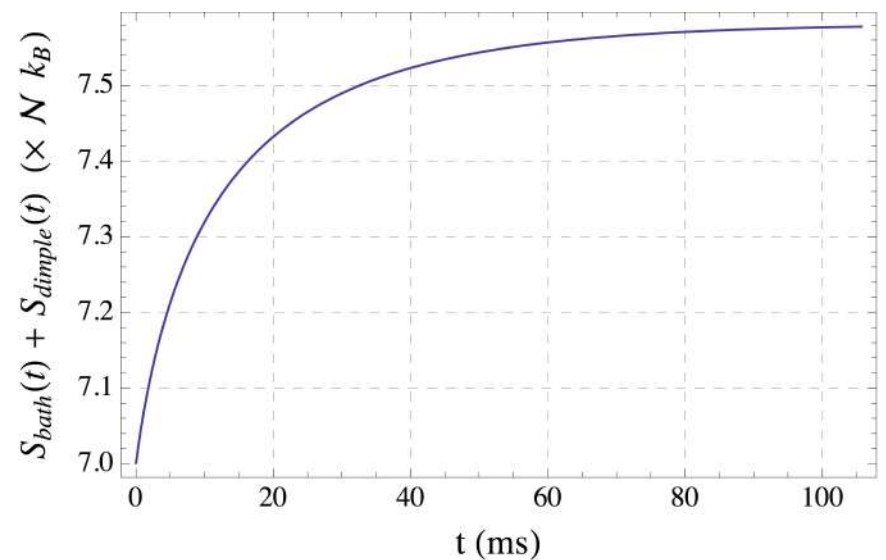
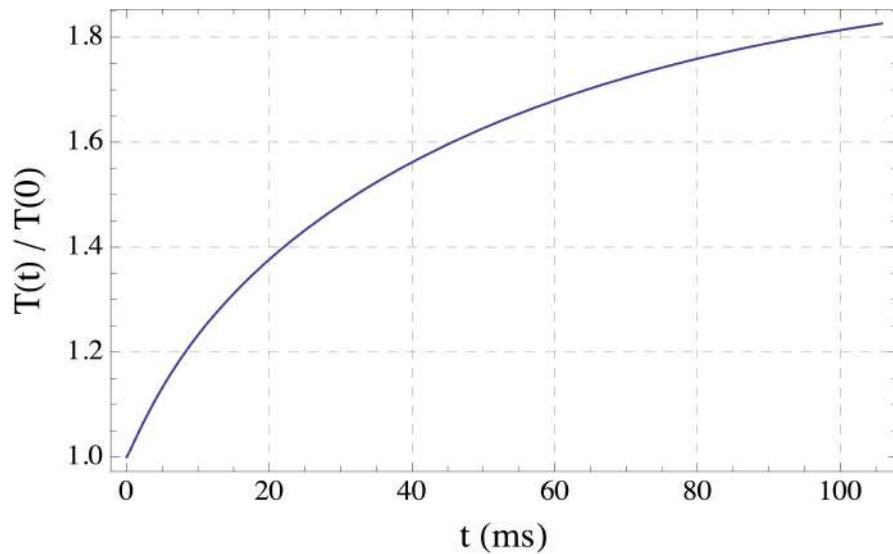
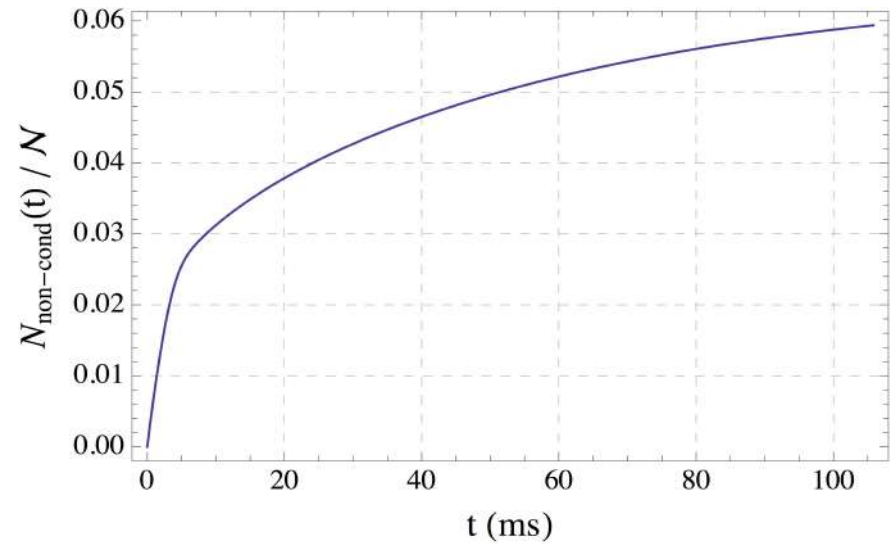
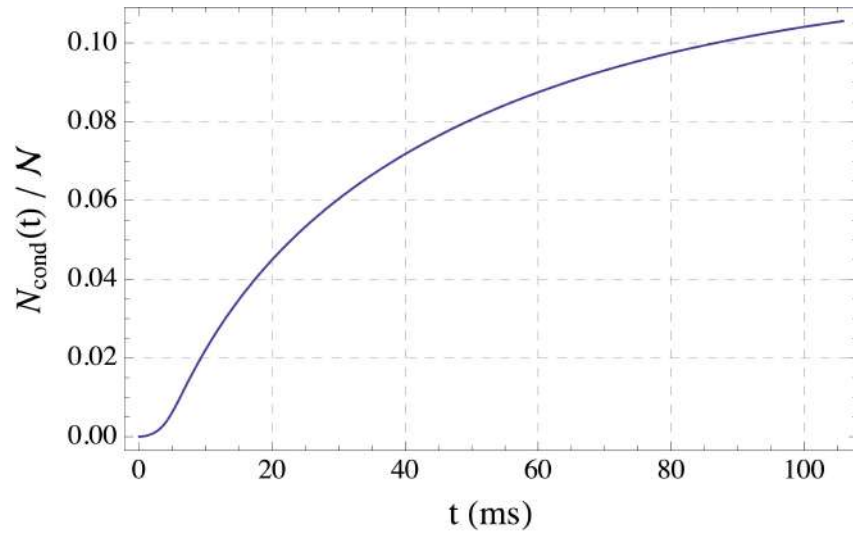


Time evolution of physical quantities

$$\beta(0) \varepsilon_{\text{dimple}} = 10$$

$$\rho_{\text{bath}}(0) \lambda_T^3(0) \approx 0.05$$

$$V_{\text{bath}}(0)/V_{\text{dimple}} = 2000$$



Thermal approximation for dimple

- ★ The dimple very quickly reaches thermal equilibrium due to interactions.
- ★ \implies Simplified model: assume a thermal distribution

$$f(\beta(0)E, t) = \frac{V_{\text{dimple}}}{\mathcal{N} \lambda_T^3(0)} \frac{2\sqrt{\beta(0)E/\pi}}{\exp[\beta_d(t)(E - \mu_d(t))] - 1}$$

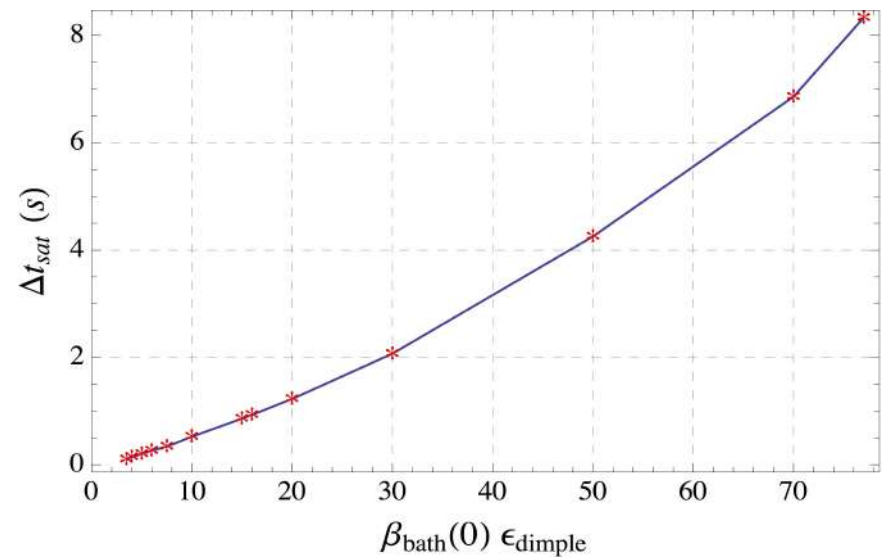
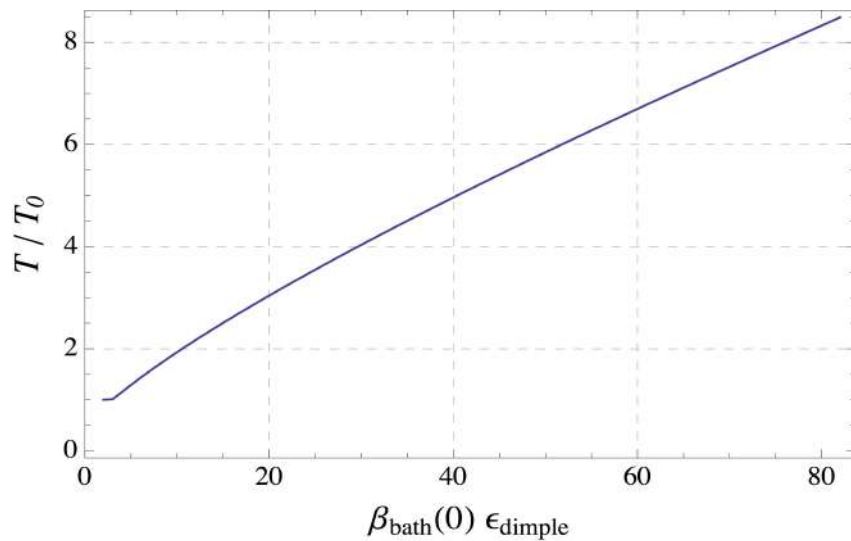
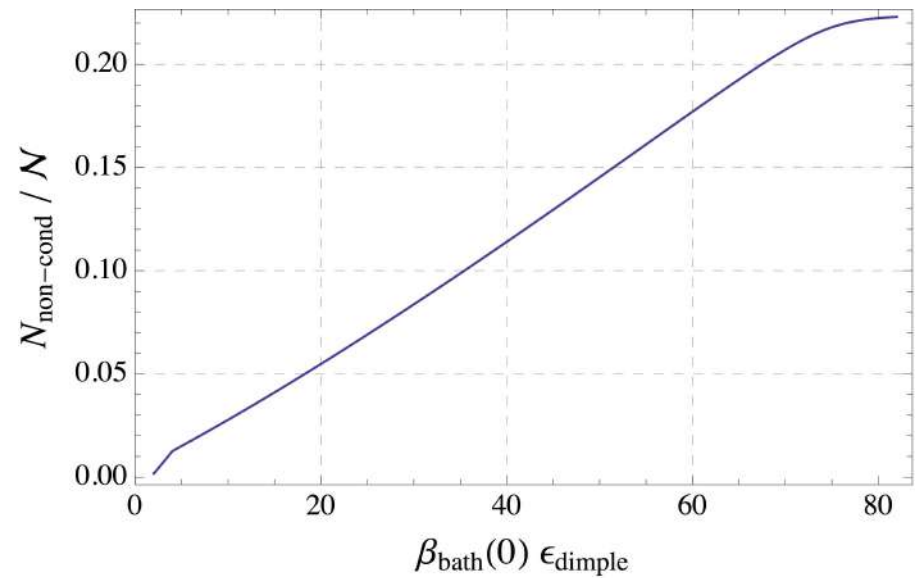
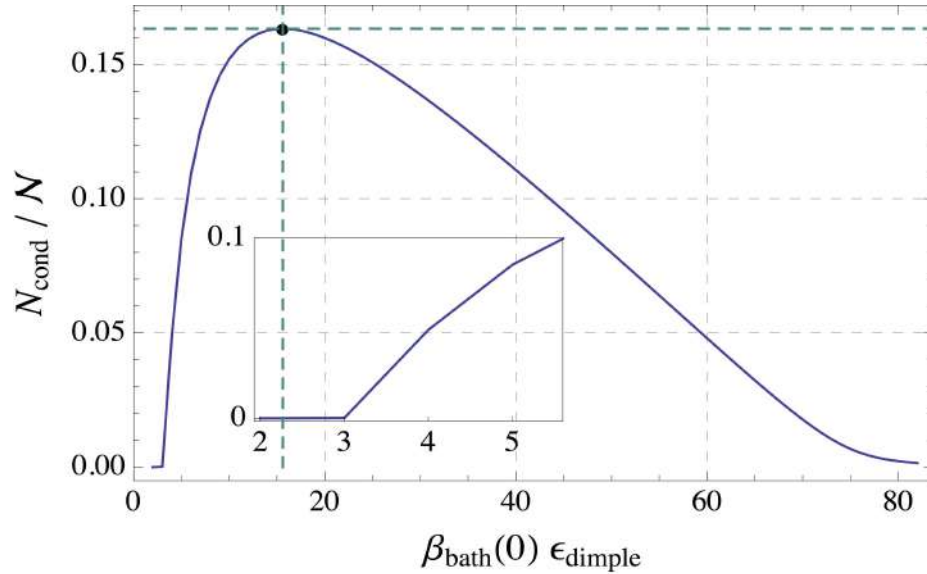
and see how $\beta_d(t)$ and $\mu_d(t)$ vary with time.

- ★ This method reproduces all features of the general model following thermal equilibrium.

Variation with dimple depth

$$\rho_{\text{bath}}(0)\lambda_T^3(0) = e^{-3}$$

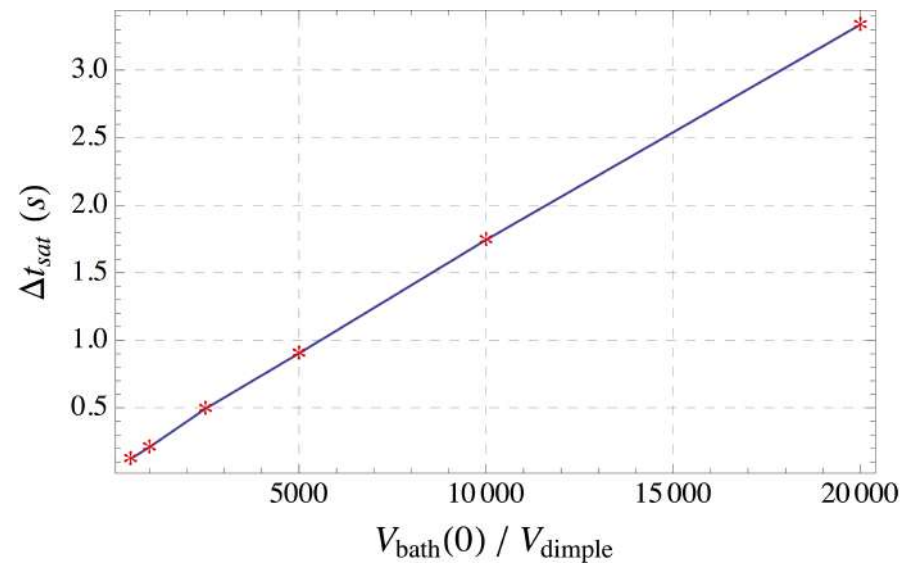
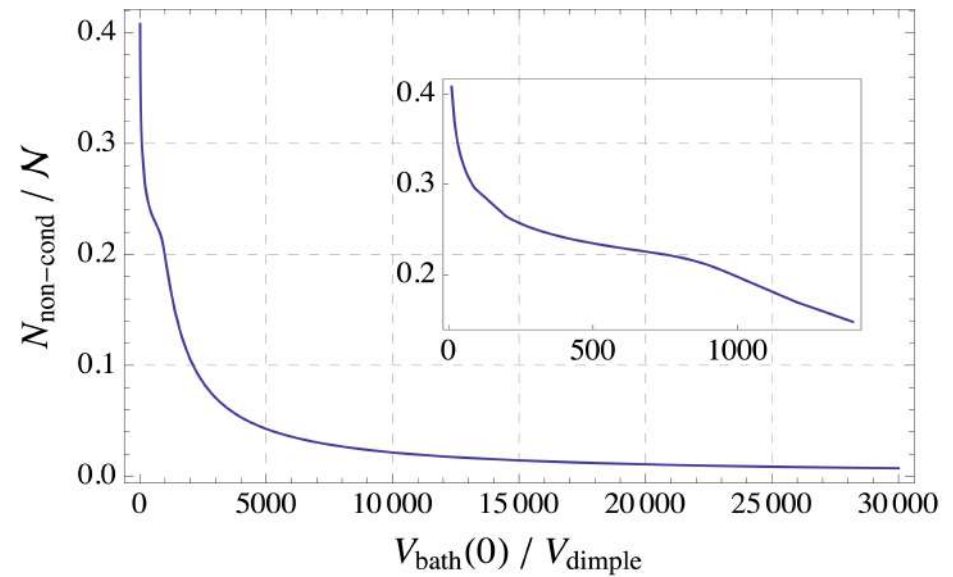
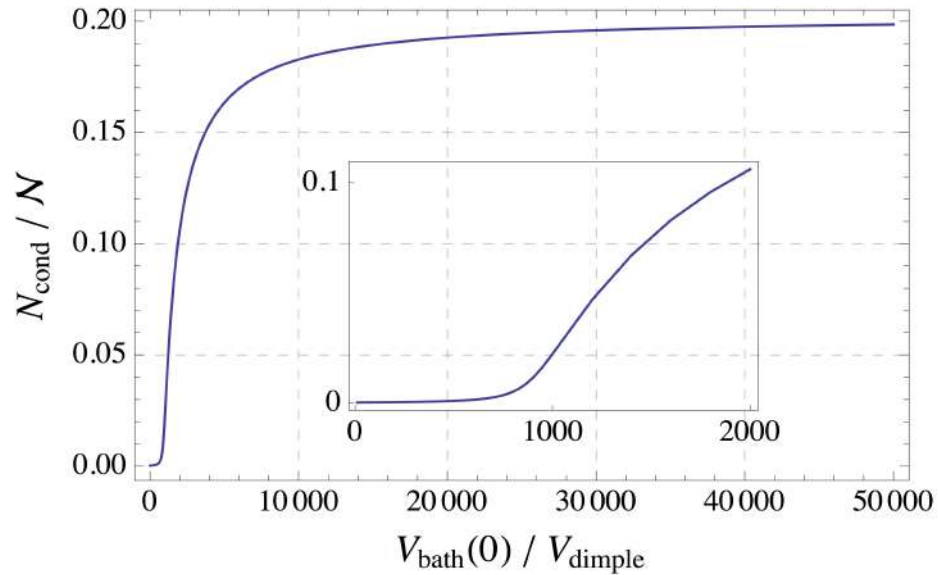
$$V_{\text{bath}}(0)/V_{\text{dimple}} = 5000$$



Variation with volume ratio

$$\rho_{\text{bath}}(0)\lambda_T^3(0) = e^{-3}$$

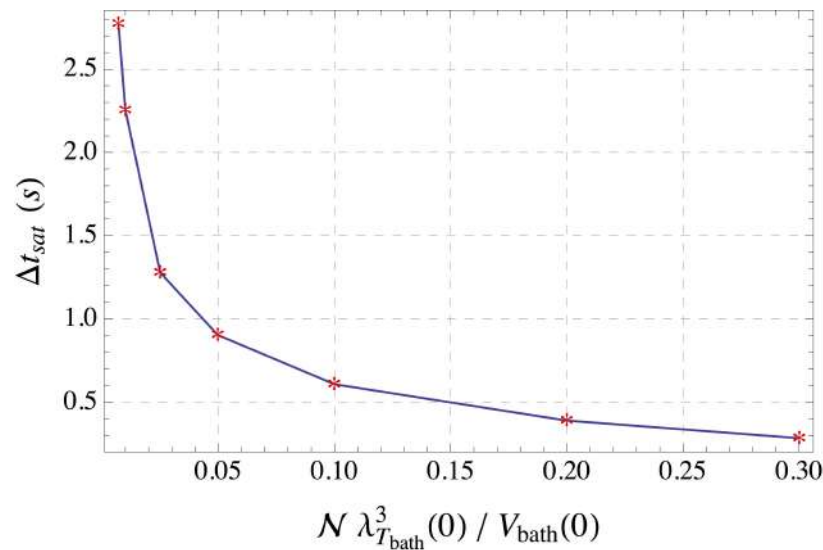
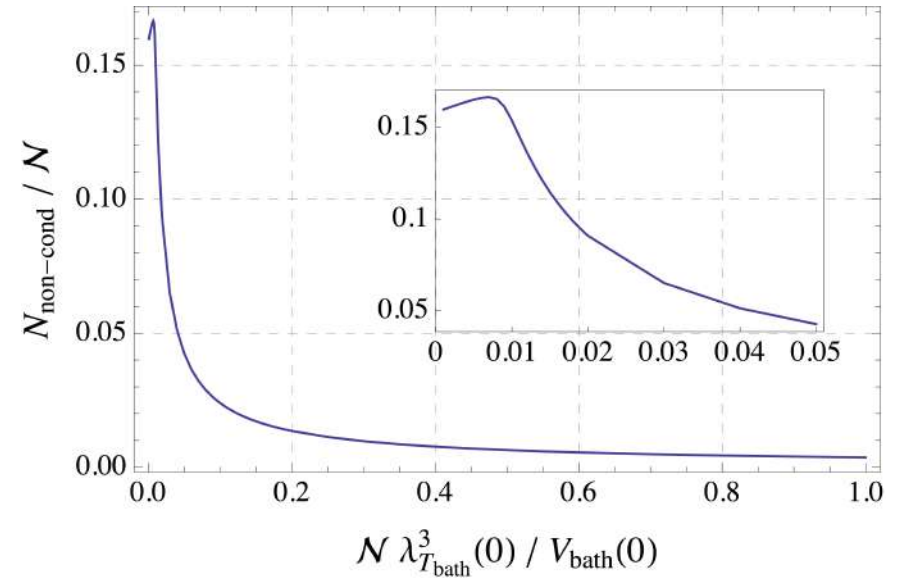
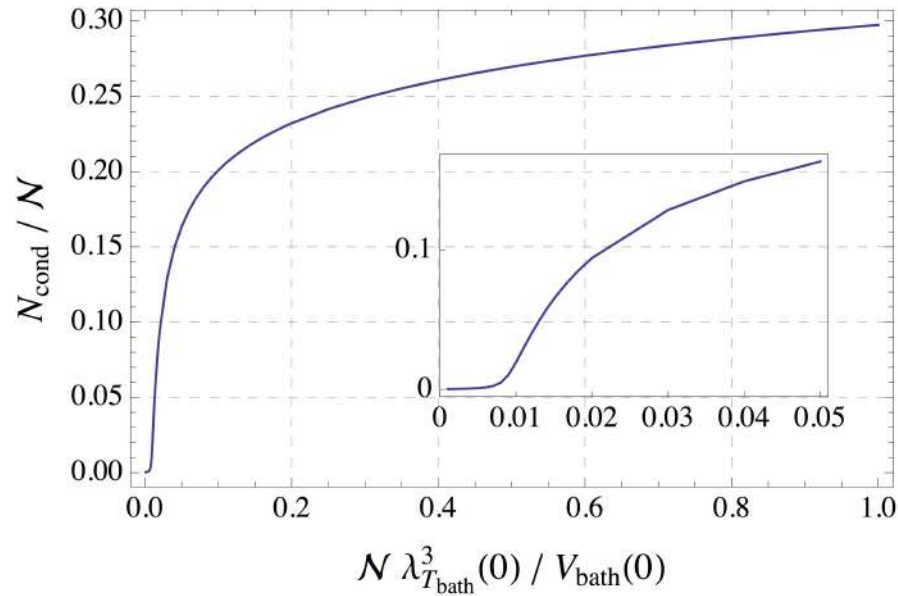
$$\beta(0)\varepsilon_{\text{dimple}} \approx 15.6$$



Variation with phase space density

$$V_{\text{bath}}(0)/V_{\text{dimple}} = 5000$$

$$\beta(0) \varepsilon_{\text{dimple}} \approx 15.6$$



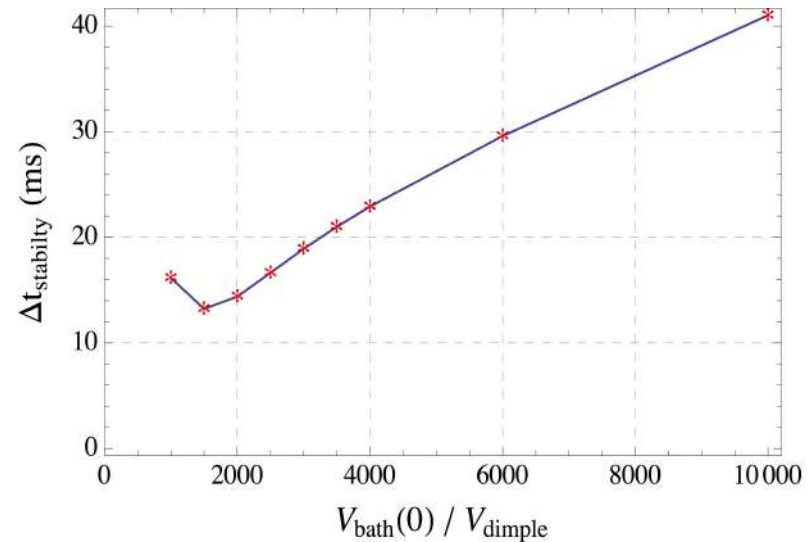
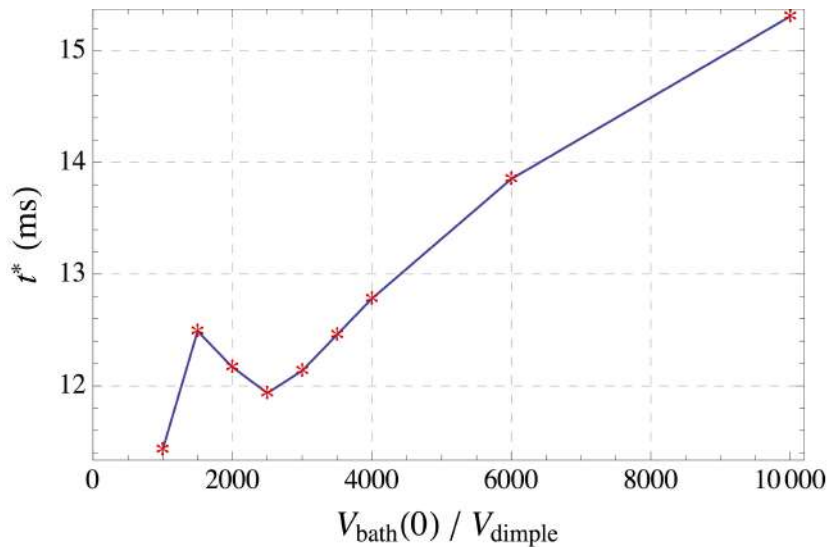
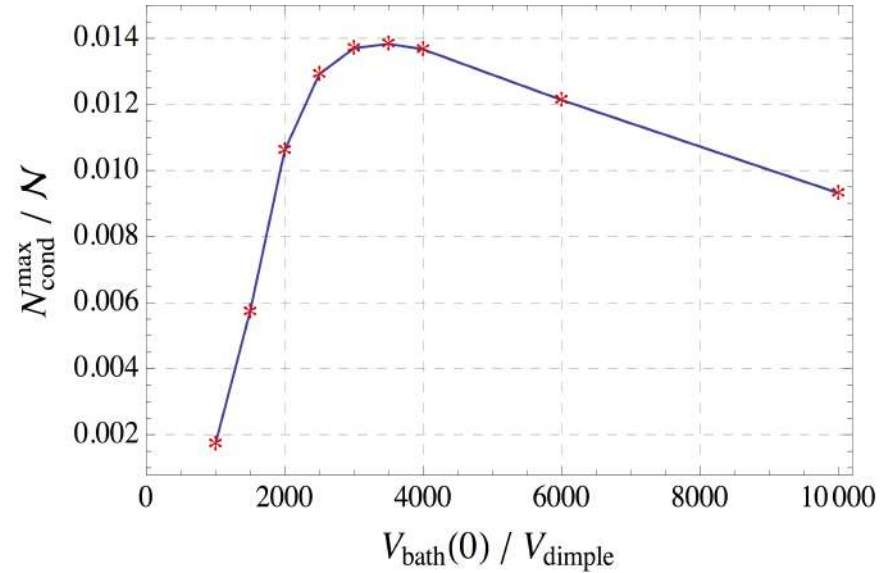
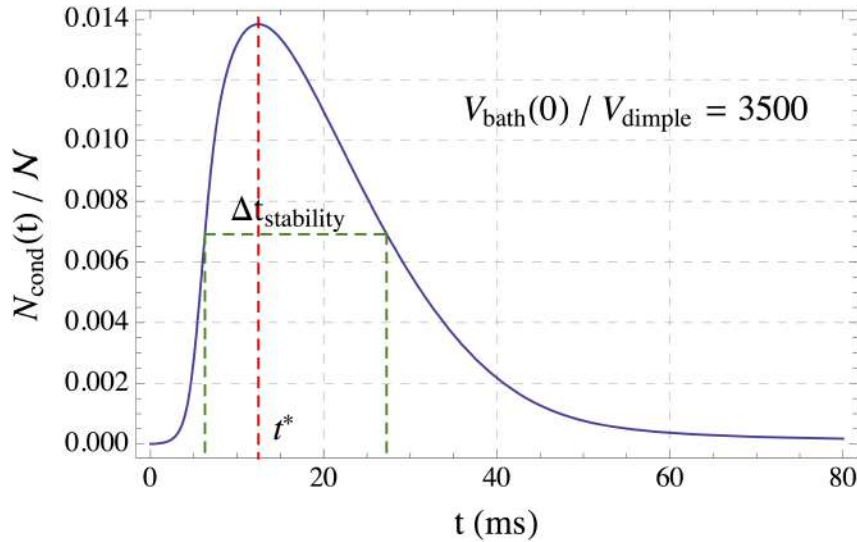
Three-body loss

- ★ A collision of 3 Rb atoms can create a Rb_2^* molecule in excited vibrational state.
- ★ The released energy causes both the atom and the molecule to escape.
- ★ Rate of such processes at position \vec{r}
$$\propto n_0^3(\vec{r}) + 9 n_0^2(\vec{r}) n_{\text{th}}(\vec{r}) + 18 n_0(\vec{r}) n_{\text{th}}^2(\vec{r}) + 6 n_{\text{th}}^3(\vec{r})$$
where $n_0(\vec{r})$ and $n_{\text{th}}(\vec{r})$ density of condensed and “thermal” particles.
- ★ This contributes terms like $-\gamma n^2(\vec{r})$, ($\gamma \approx 10^{-29} \text{cm}^6/\text{s}$) which destabilize condensates at high density.

Effect of three-body loss on condensate fraction

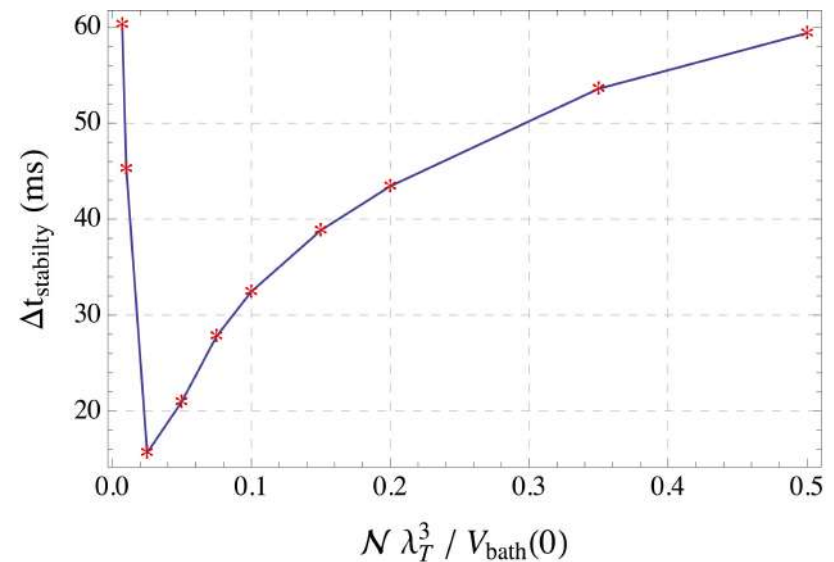
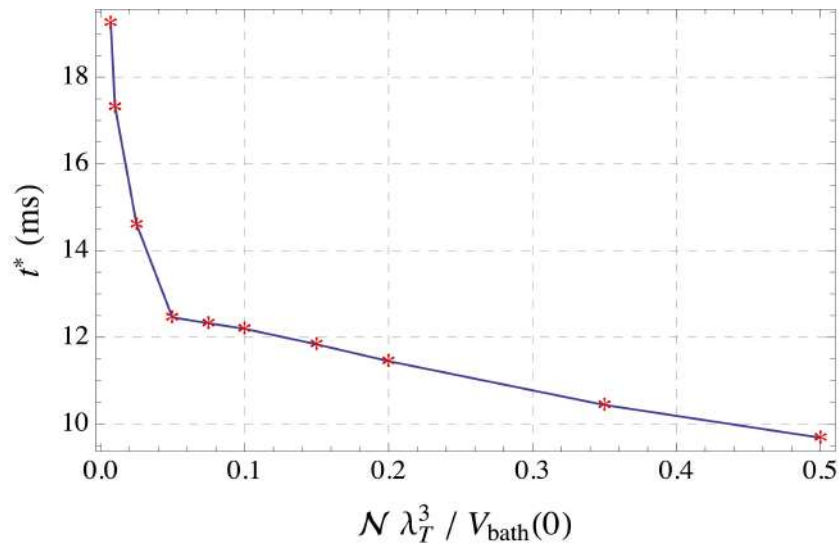
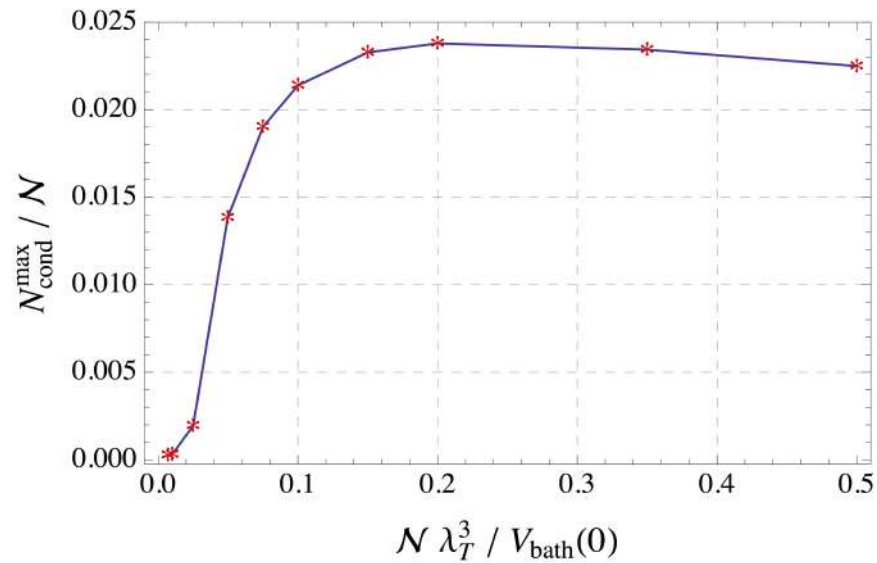
$$\rho_{\text{bath}}(0)\lambda_T^3(0) = e^{-3}$$

$$\beta(0)\varepsilon_{\text{dimple}} \approx 15.6$$

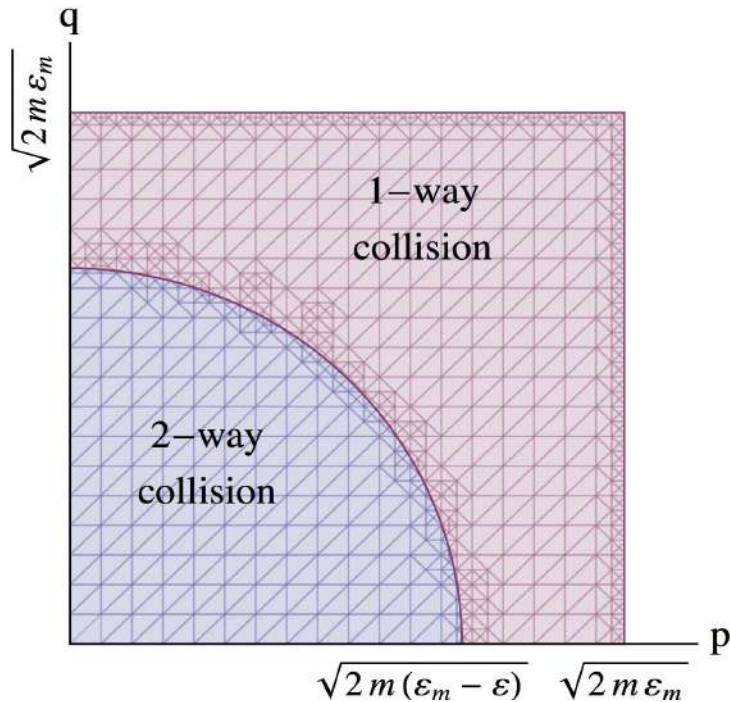


Effect of three-body loss on condensate fraction

$$V_{\text{bath}}(0)/V_{\text{dimple}} = 3500$$
$$\beta(0) \varepsilon_{\text{dimple}} \approx 15.6$$



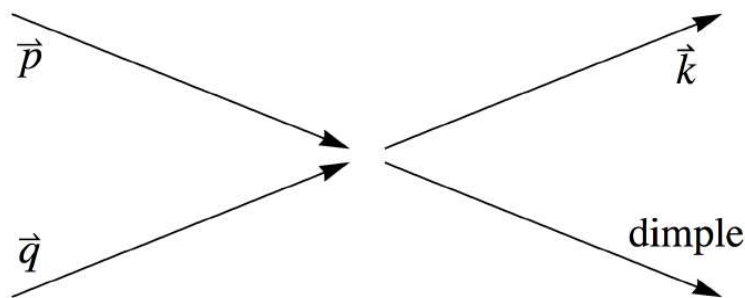
Effect of finite trap depth



- ★ 2-way collisions: no particle is lost
- ★ 1-way collisions: one particle goes to the dimple and the other is lost
 \Rightarrow energy lost from bath = $\frac{(p^2 + q^2)}{2m}$

- ★ This causes cooling, resulting in higher condensate fraction

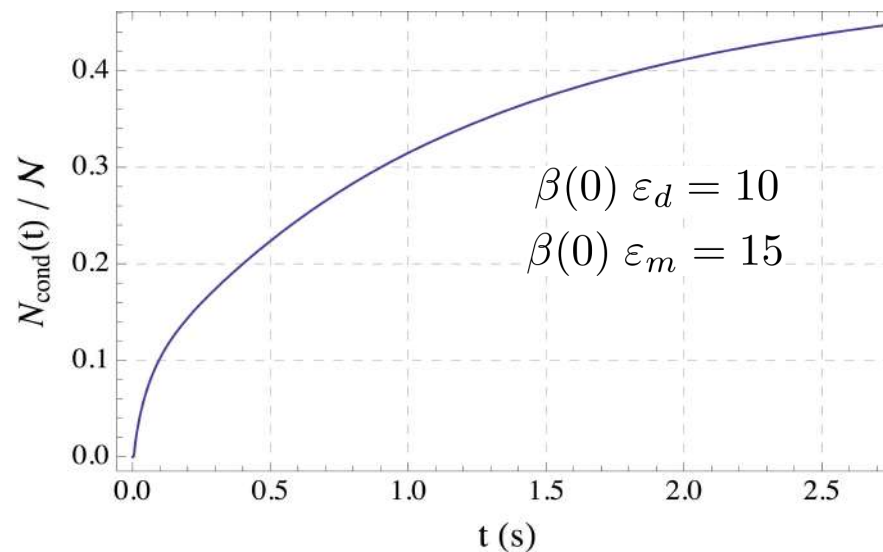
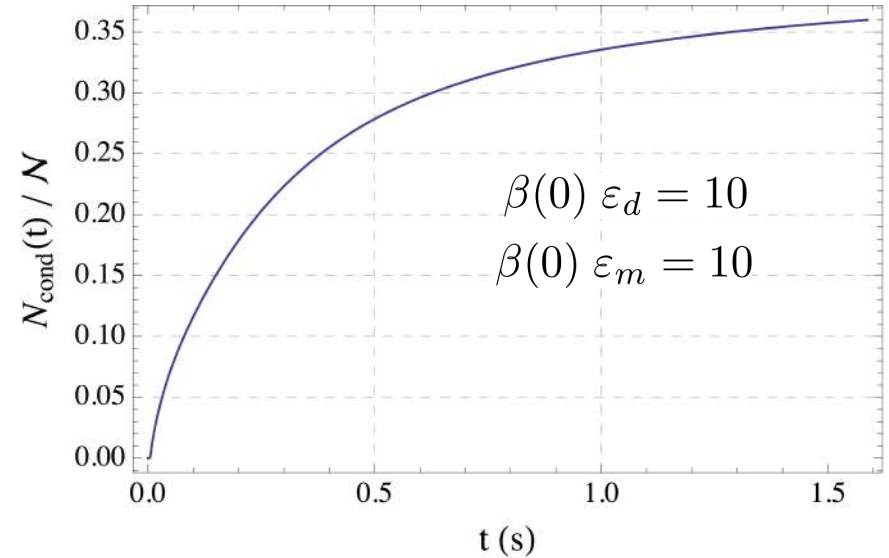
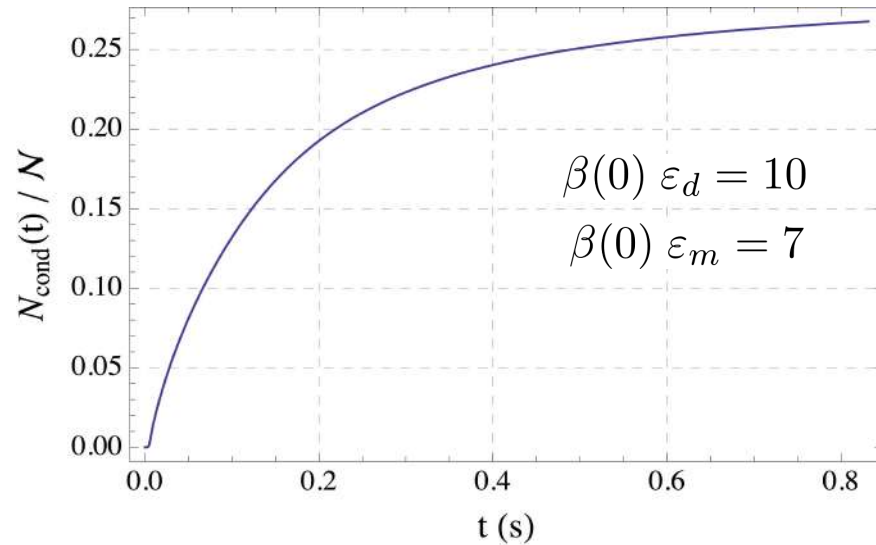
- ★ For $\epsilon_m \gg \epsilon$, 1-way collisions become negligible



Plots for finite trap depth (without 3 body loss)

$$\rho_{\text{bath}}(0)\lambda_T^3(0) = e^{-3}$$

$$V_{\text{bath}}(0)/V_{\text{dimple}} = 5000$$



Thank you! 😊

Questions?

Shovan Dutta and Erich J. Mueller
arXiv:1407.2557