

# Pulsed Generation of Symmetry-Protected Long-Range Entanglement

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SD, Stefan Kuhr, & Nigel Cooper, arXiv 2201.10564

**EPSRC**

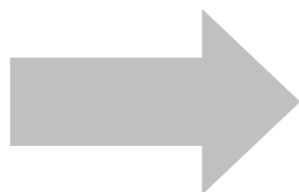
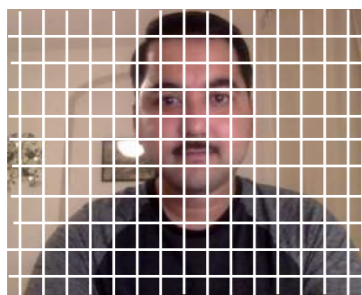
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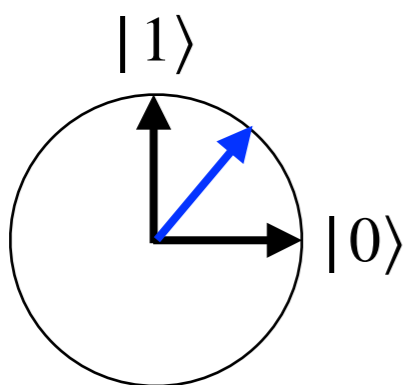
# Qubits and Entanglement

Classical bit — unit of information: 0 or 1



Data stream ...0010110101...

Quantum bit (“qubit”) — can be in a **superposition**  $\alpha|0\rangle + \beta|1\rangle$



Measurement gives:

$|0\rangle$  : probability  $|\alpha|^2$

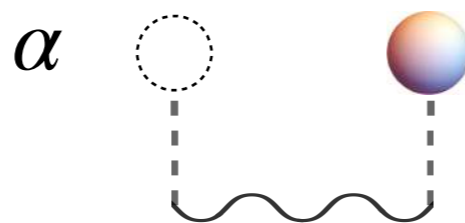
$|1\rangle$  : probability  $|\beta|^2$

Spin:  $\alpha|\downarrow\rangle + \beta|\uparrow\rangle$

Lattice site:  $\alpha|\circ\rangle + \beta|\bullet\rangle$

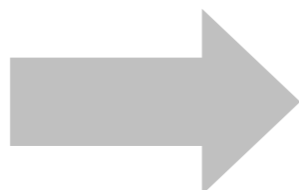
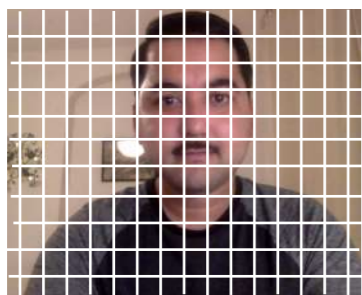
**Entangled** pair of qubits  $\alpha|01\rangle + \beta|10\rangle$

$|\alpha| = |\beta|$  : “Bell pair”



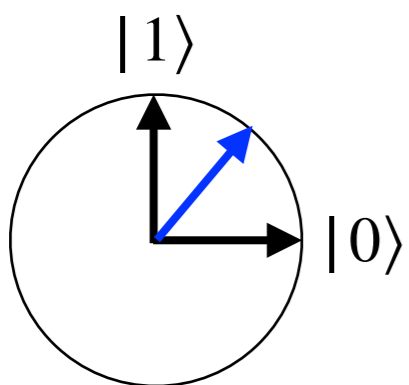
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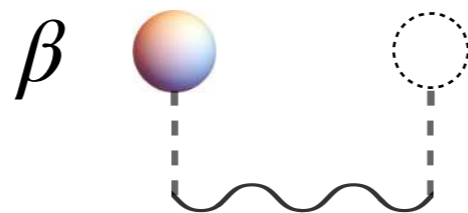
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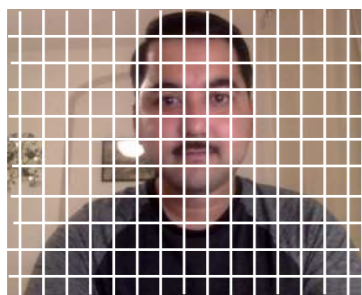
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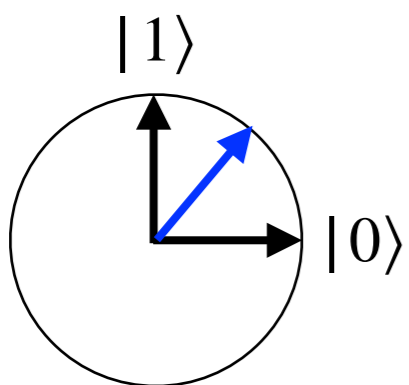
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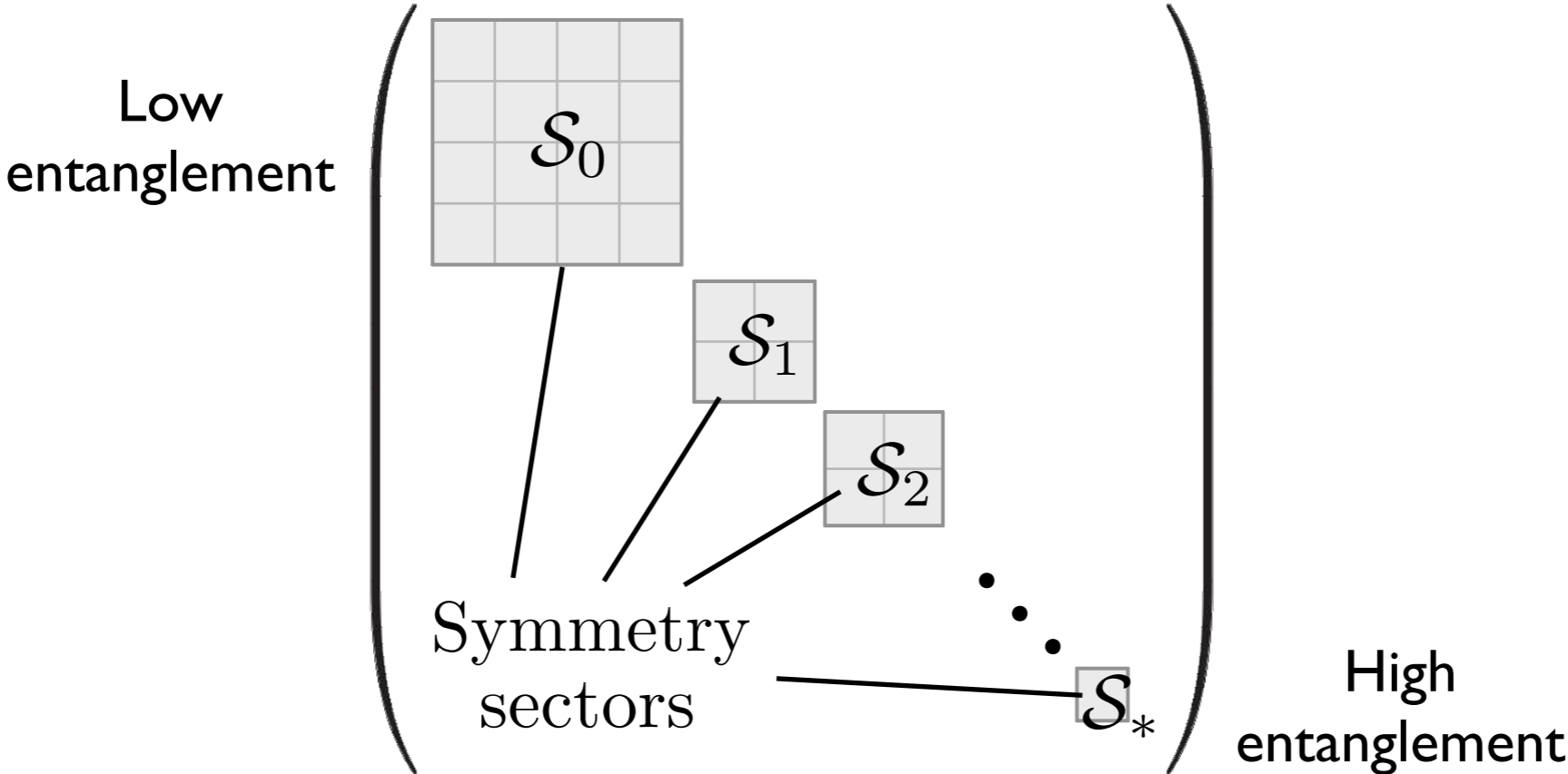
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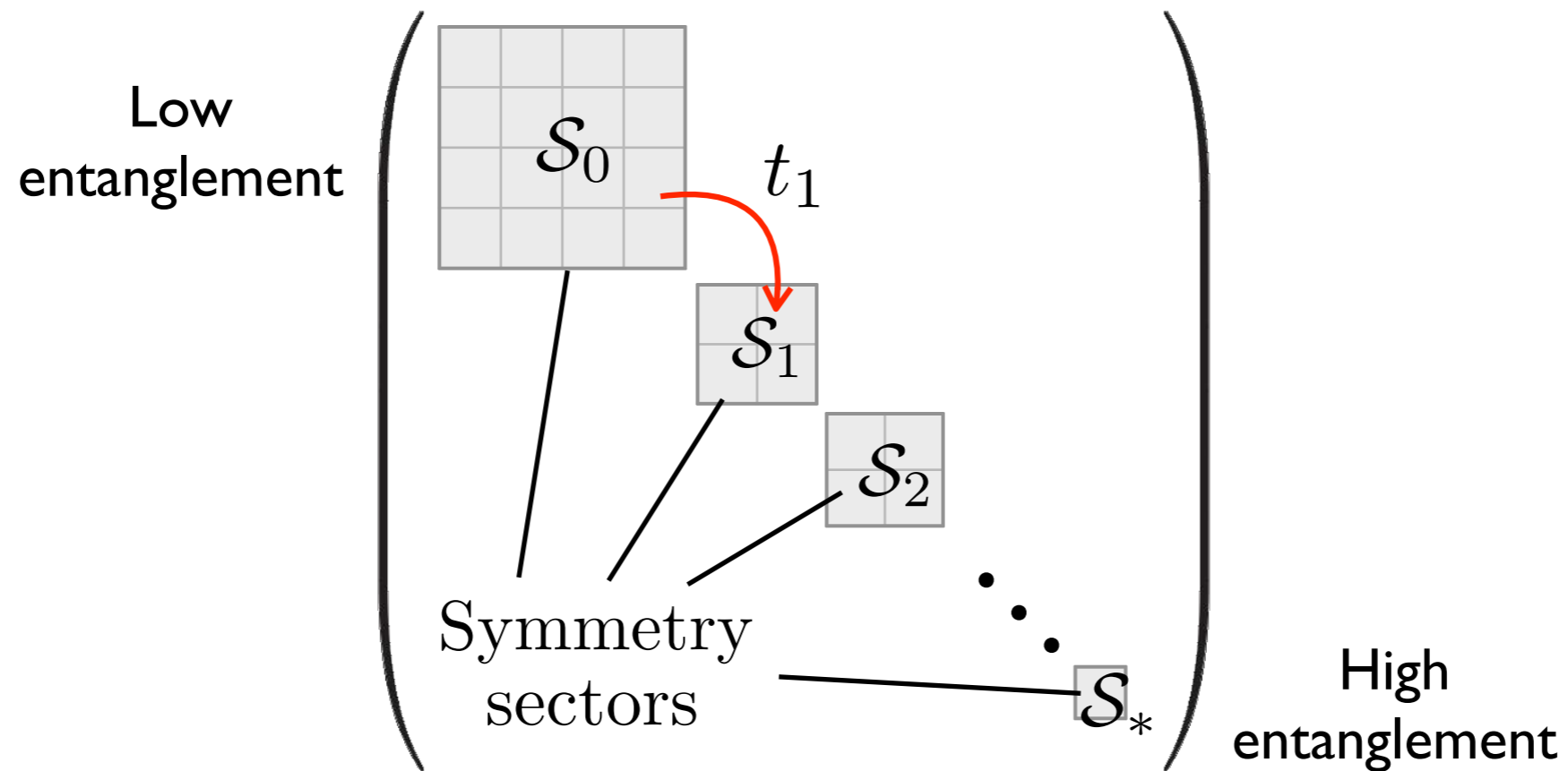


Hard to prepare in a generic many-body environment

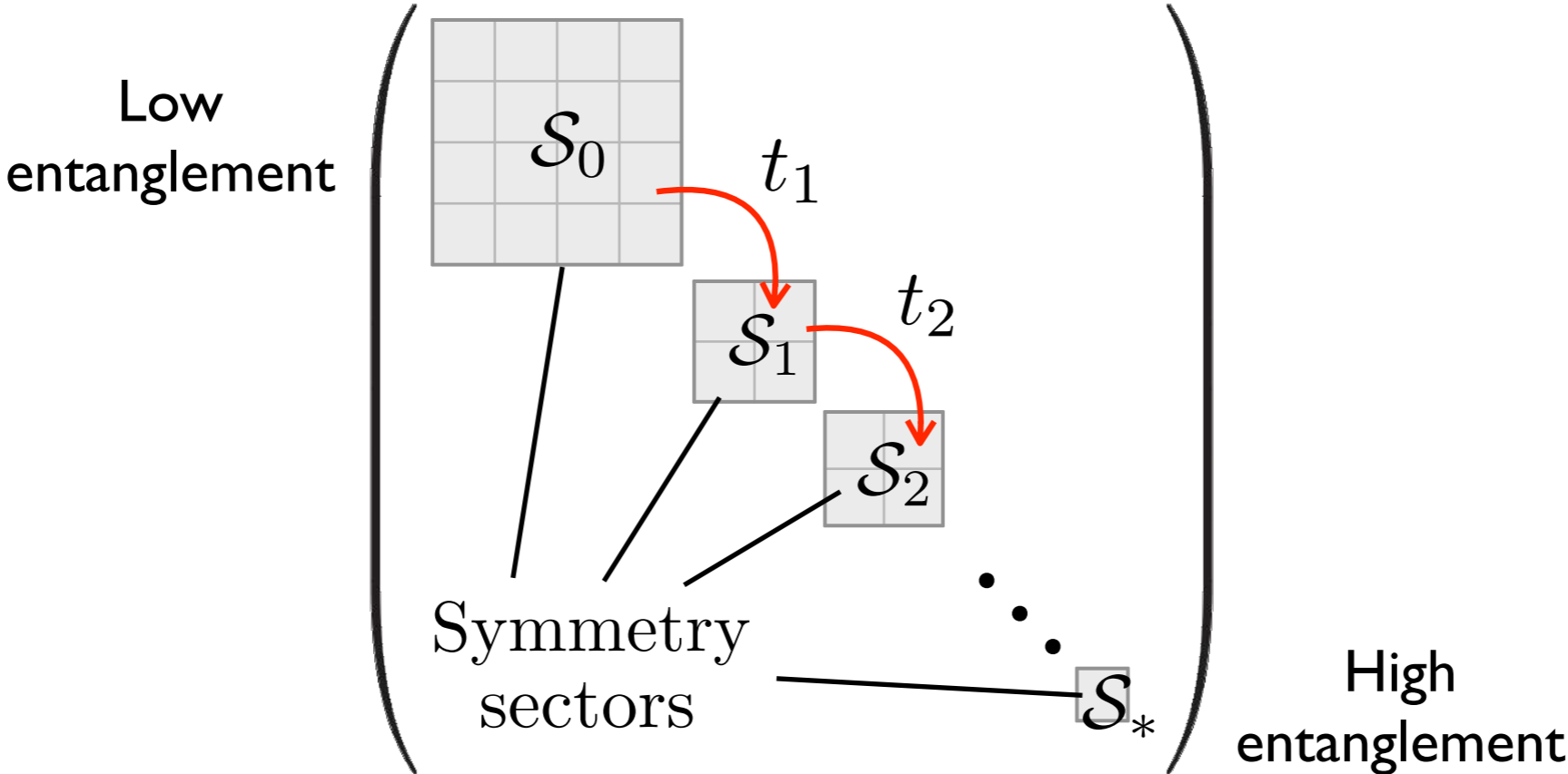
## Harness symmetry to generate long-range entanglement



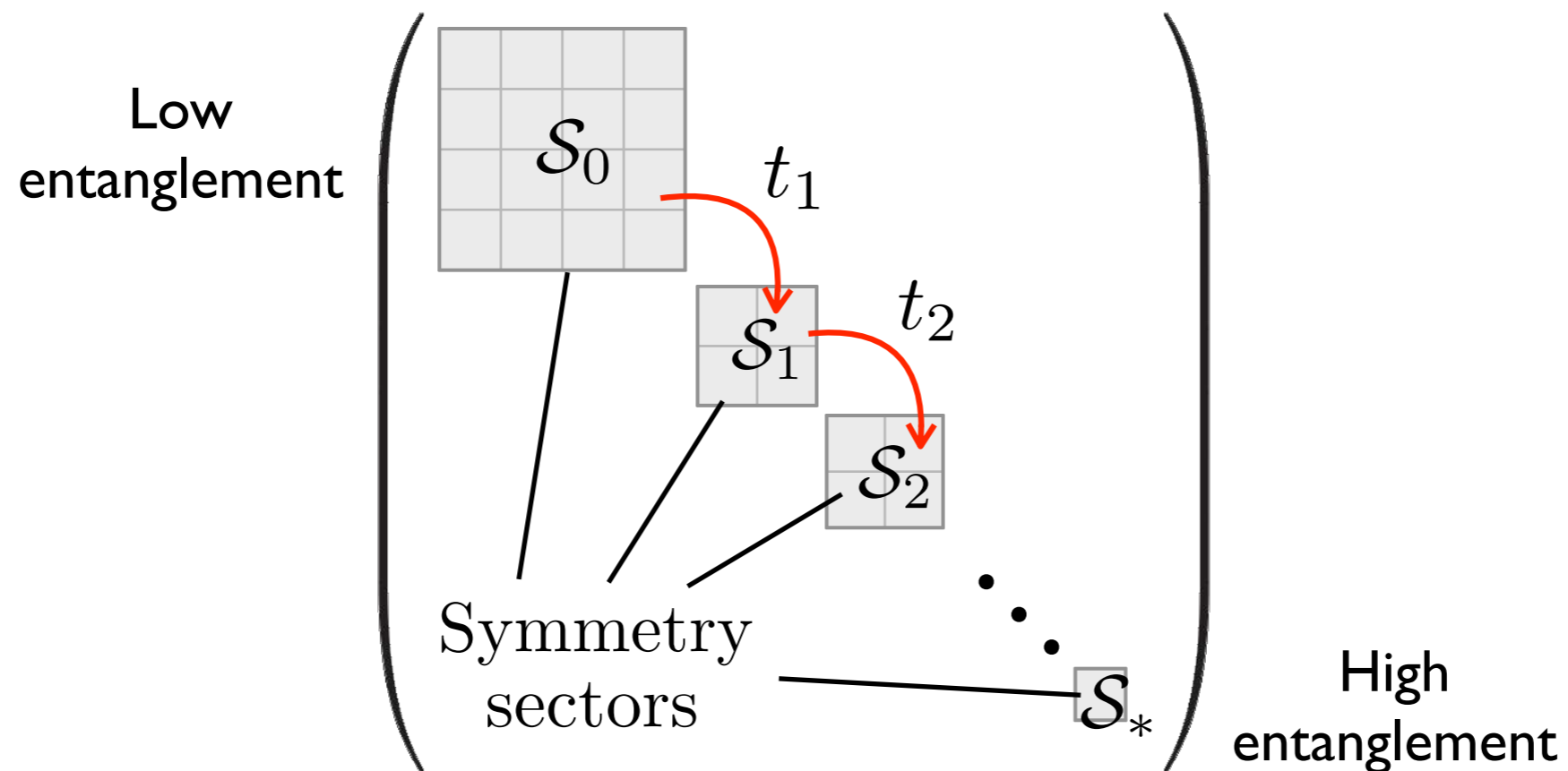
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⇒ On-demand generation of entanglement and quantum correlations



# Qubit array

Spin-1/2 XX chain

$$\hat{H} = -J \sum_i \hat{\sigma}_i^x \hat{\sigma}_{i+1}^x + \hat{\sigma}_i^y \hat{\sigma}_{i+1}^y$$

Cold atoms/ions,  
superconducting circuits...



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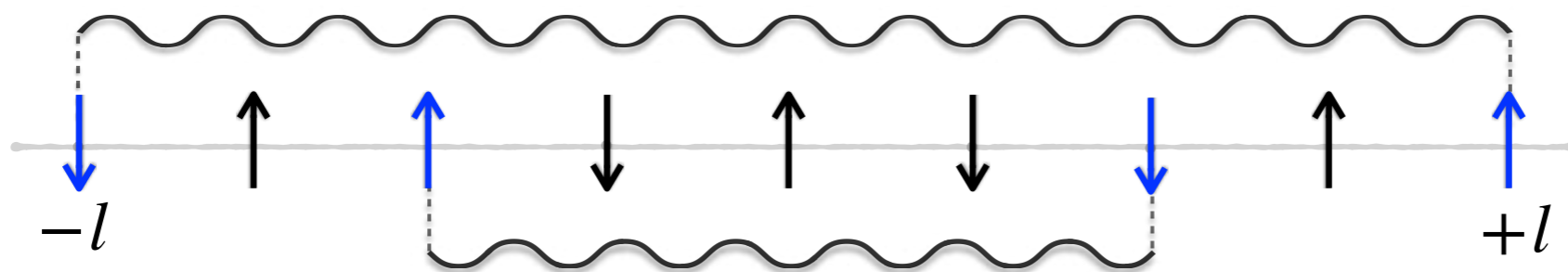
Symmetry that counts Bell pairs between sites  $\{i, -i\}$

PRL **125**, 240404 (2020)

$$[\hat{H}, \hat{C}] = 0$$

$$\hat{C} = \hat{\sigma}_0^z / 2 + \sum_{i \neq 0} \hat{f}_i^\dagger \hat{f}_{-i}$$

$$\hat{f}_i := \hat{\sigma}_i^- \prod_{j < i} \hat{\sigma}_j^z$$



Strategy: Produce  $l$  Bell pairs by pumping  $C$

# Procedure

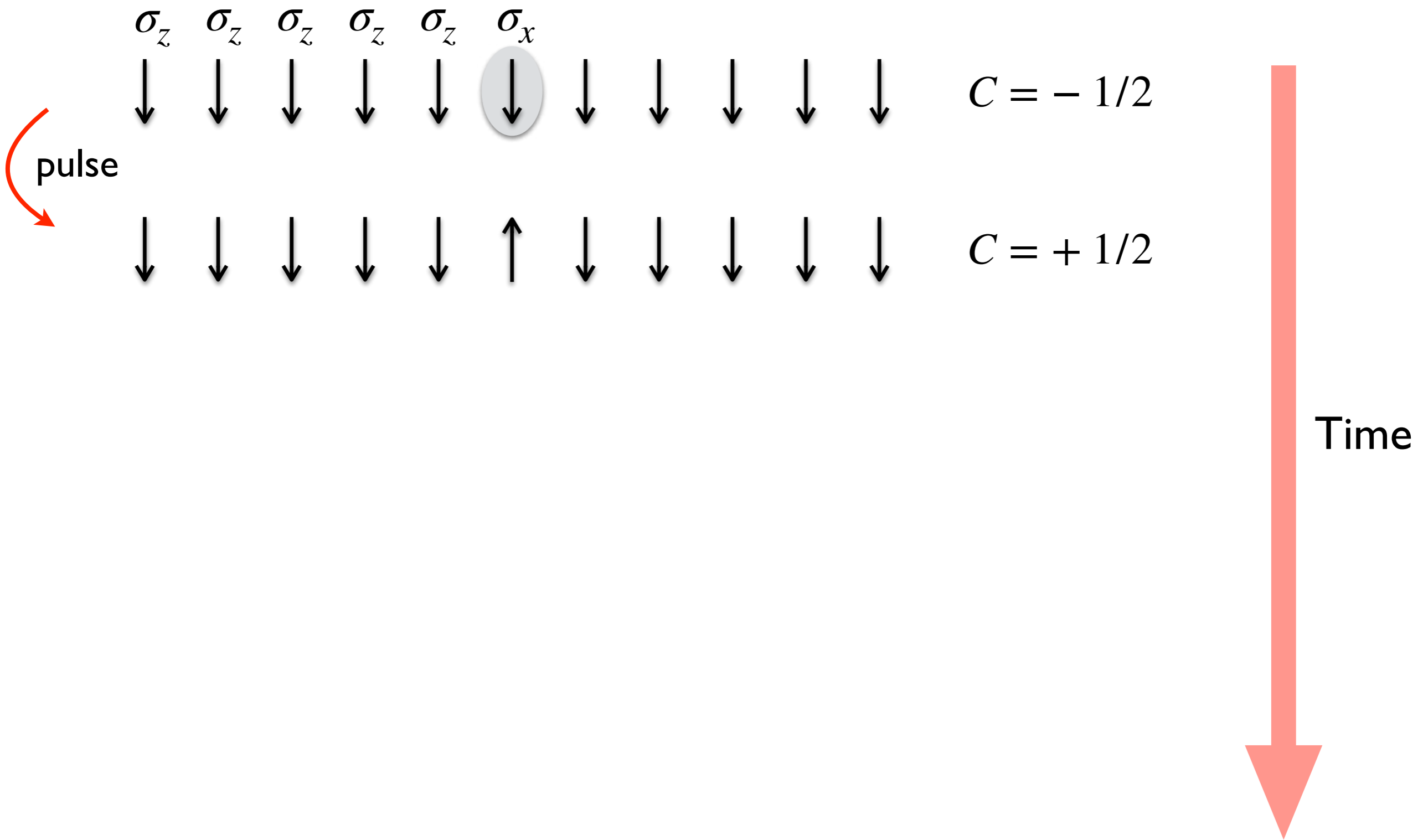


$$C = -1/2$$

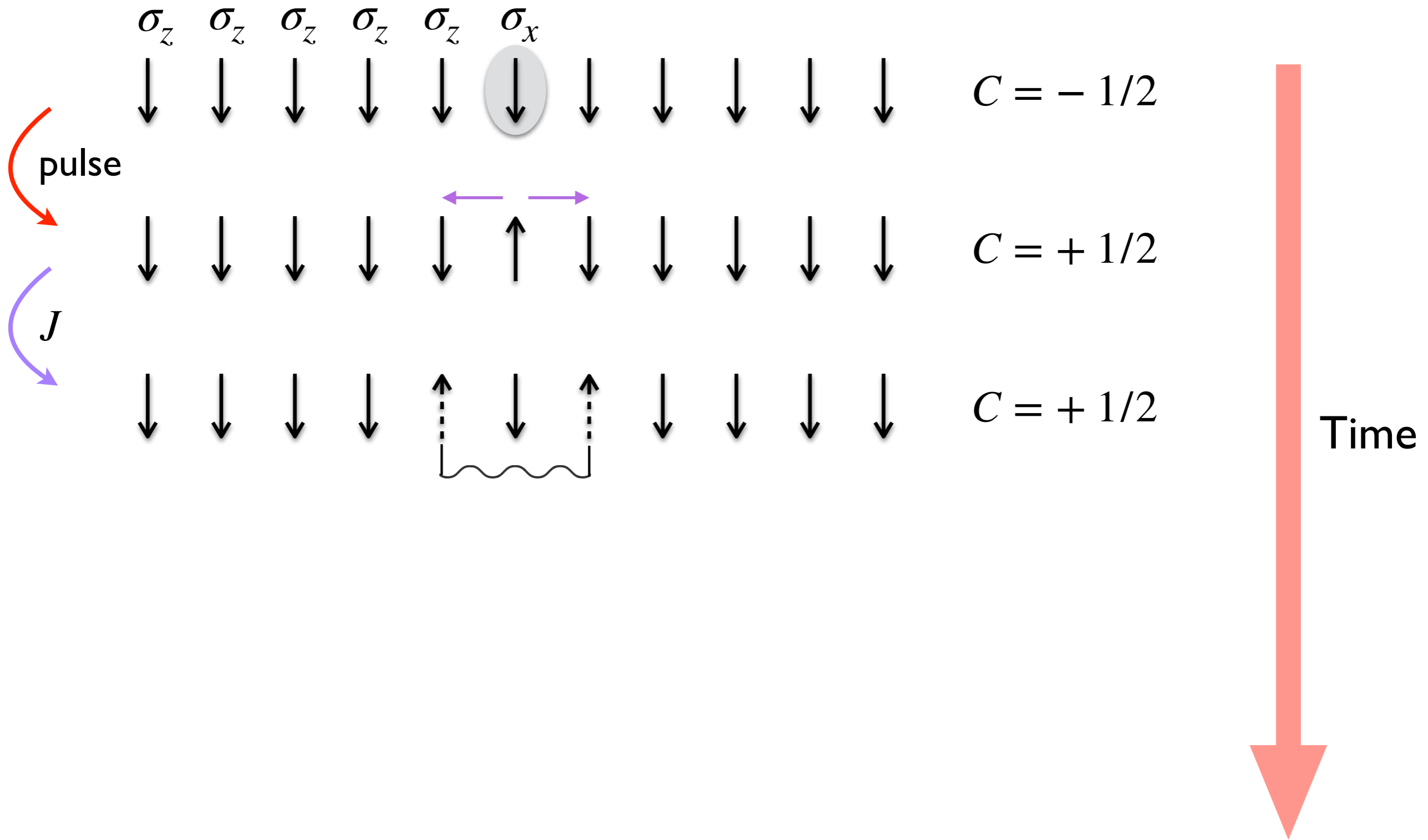


Time

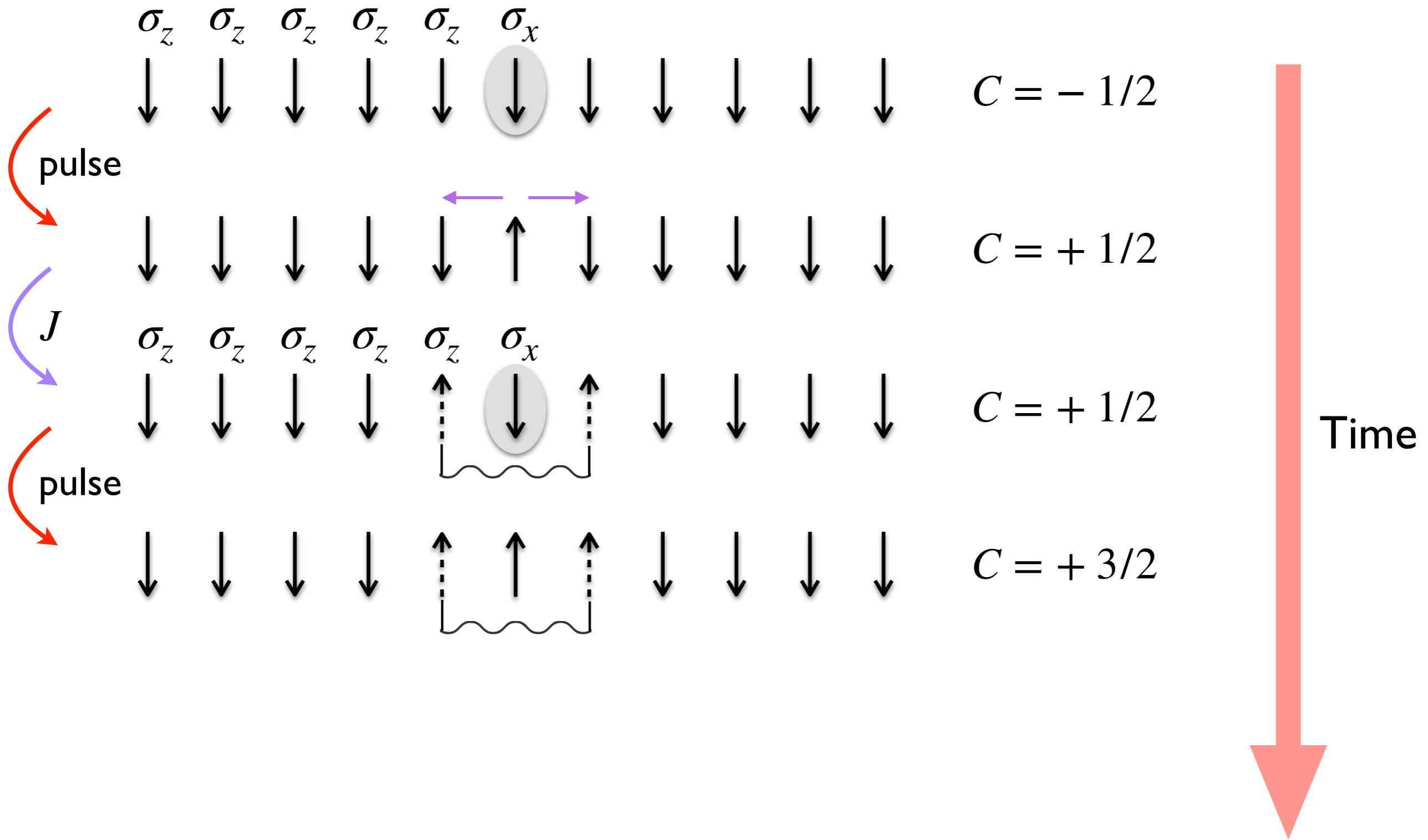
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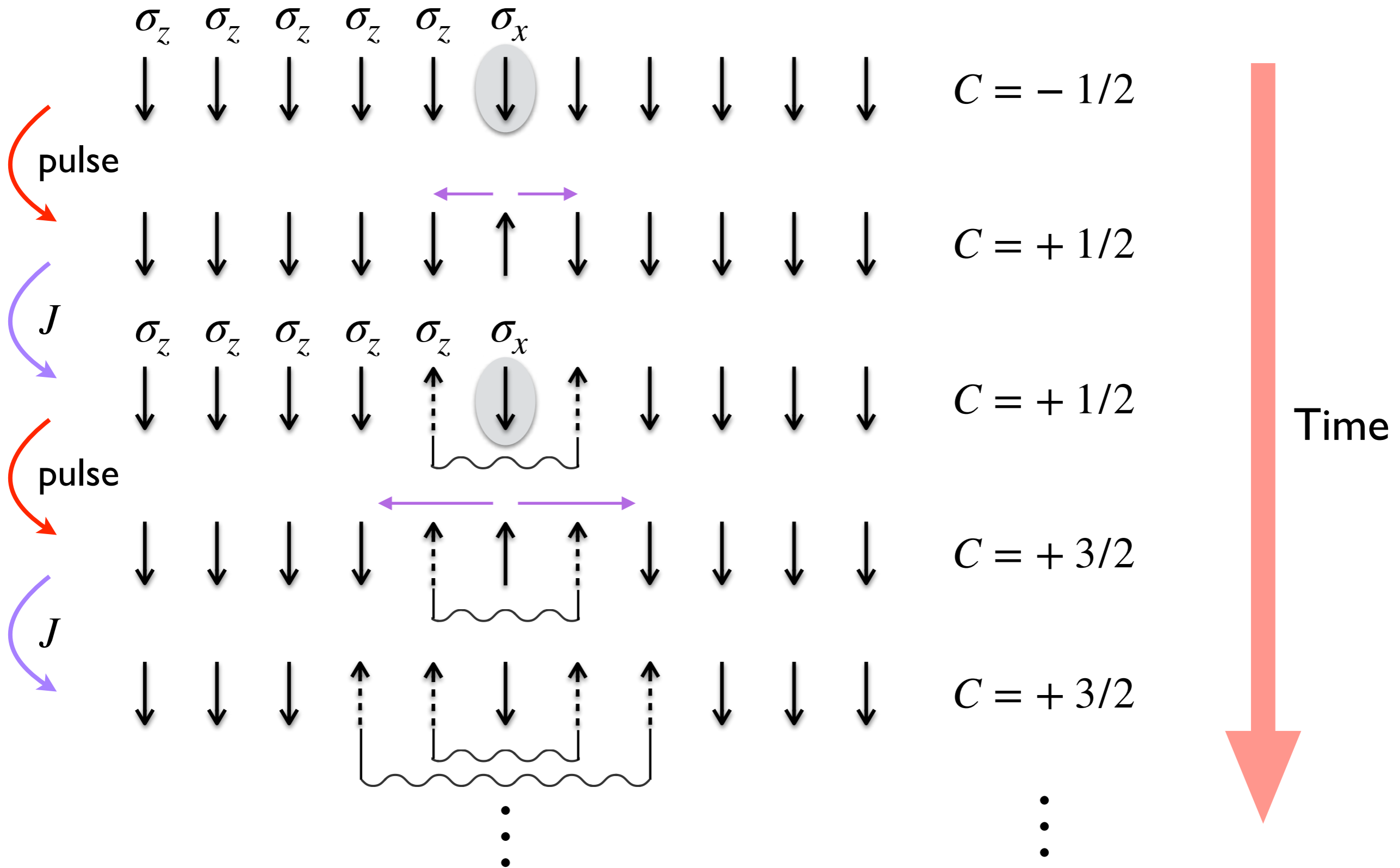
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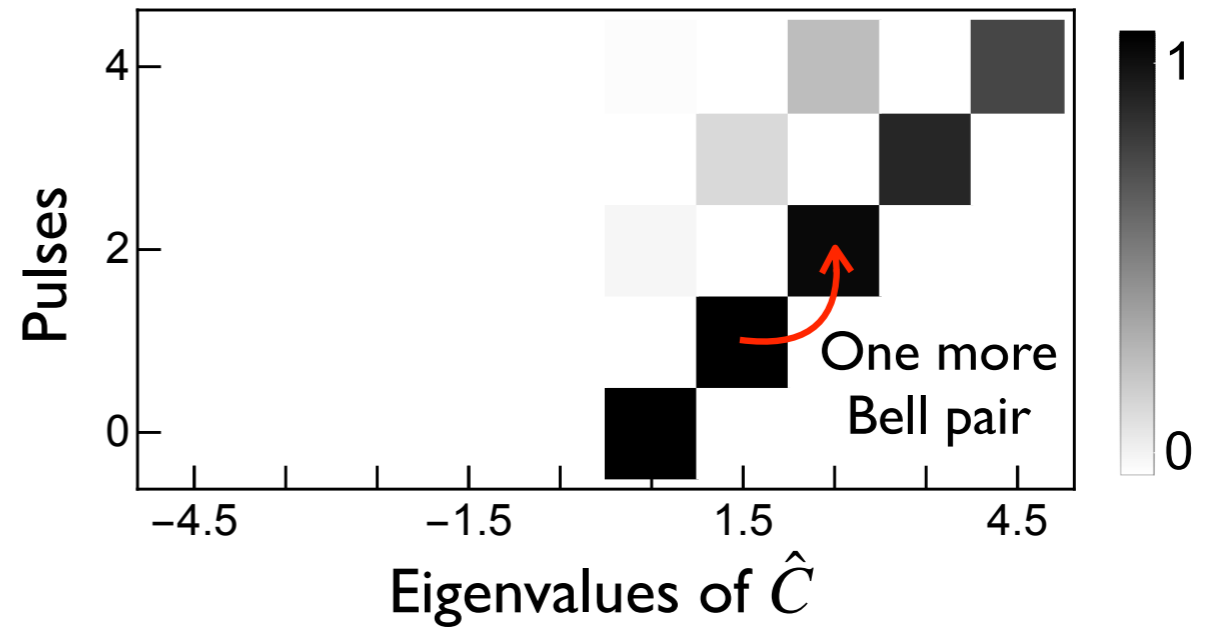
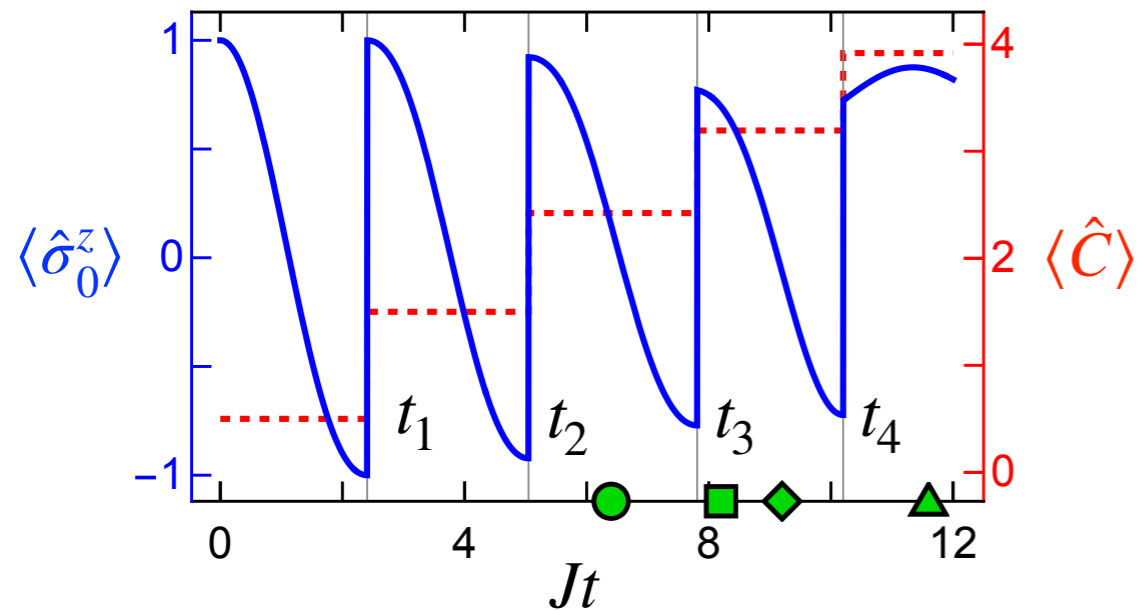
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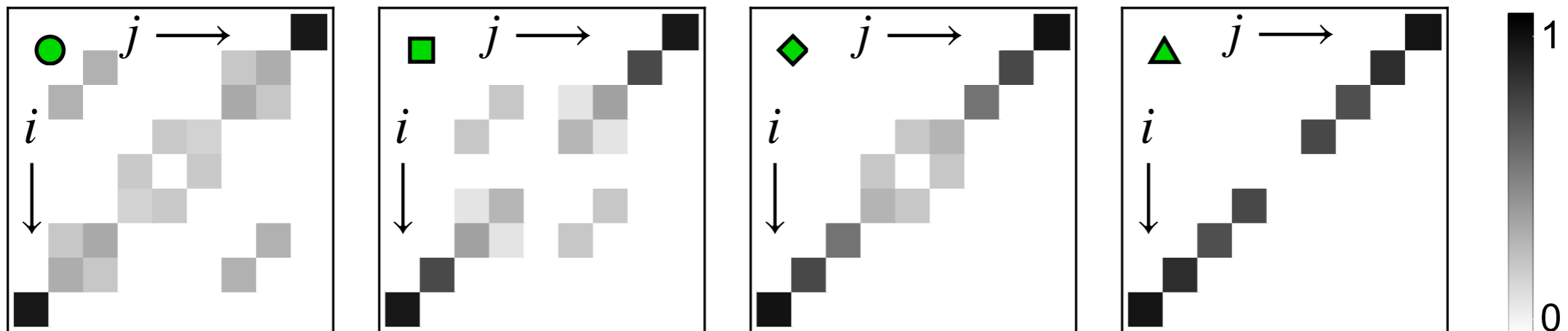


# Result — Exact diagonalisation

Pulse whenever centre site is mostly  $\downarrow$ , up to  $l$  times

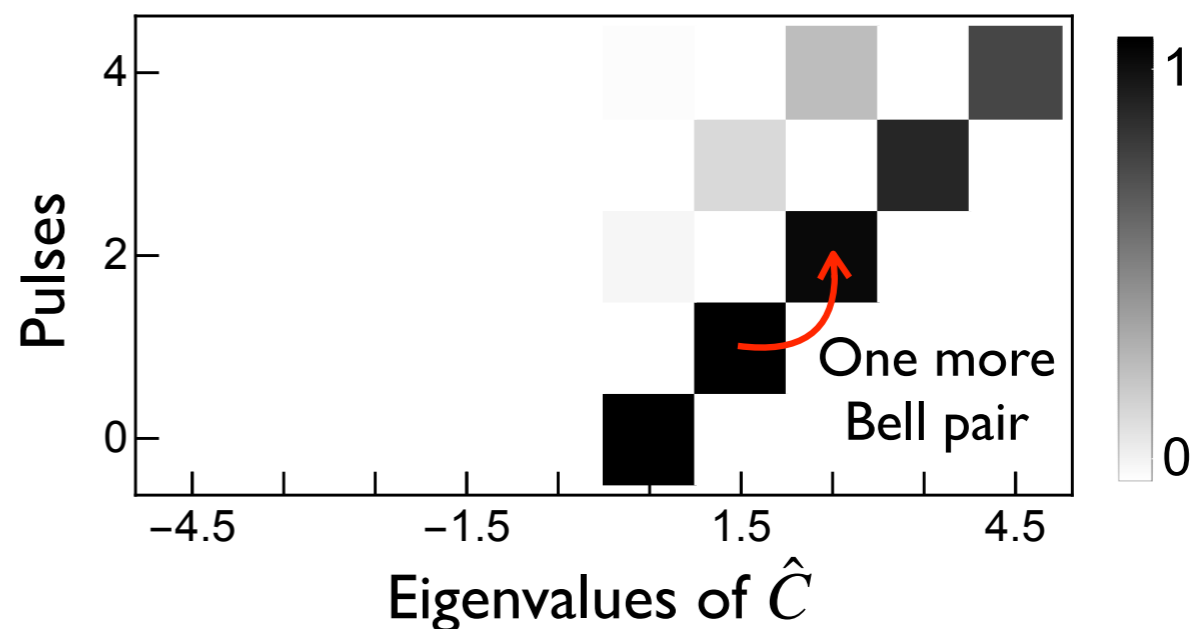
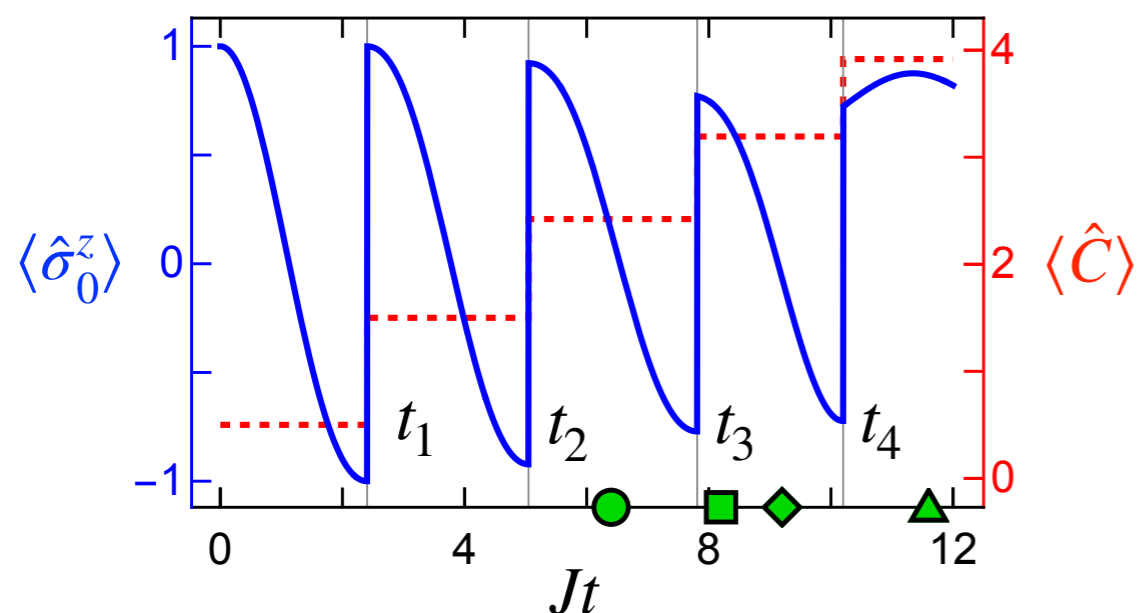


Concurrences  $\mathcal{C}_{i,j}$  show Bell pairs are successively stacked inward

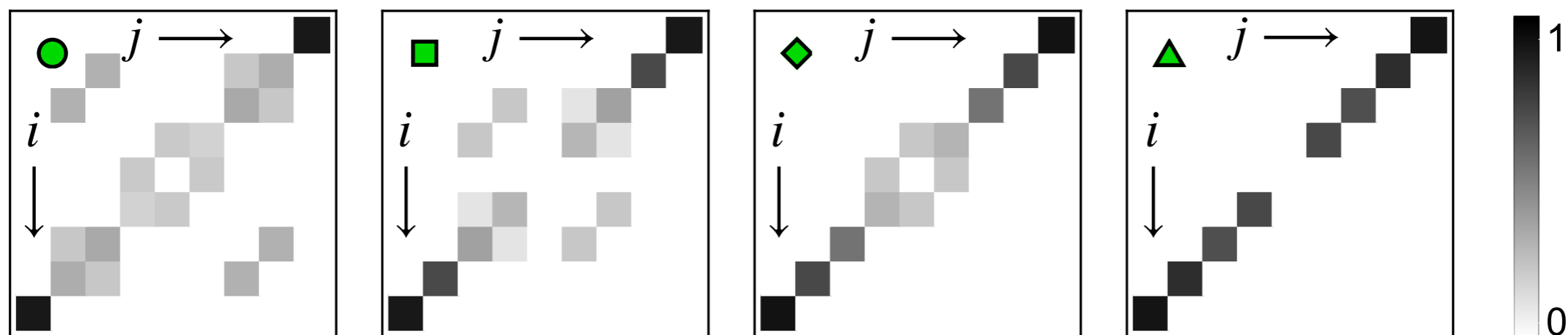


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For large  $l$  — number of Bell pairs grows linearly,  $\langle \hat{C} \rangle \sim 0.6l + 1$   
 — preparation time grows linearly,  $\Delta t \sim 2l/J$

# Fermi-Hubbard model — $\eta$ pairs

Hamiltonian  $\hat{H} = \sum_{s=\uparrow,\downarrow} \sum_{i=-l}^{l-1} (-J \hat{c}_{i,s}^\dagger \hat{c}_{i+1,s} + \mathbf{h.c.}) + U \sum_{i=-l}^l \hat{n}_{i,\uparrow} \hat{n}_{i,\downarrow}$

SU(2) symmetry:  $\hat{\eta}^- = \sum_i (-1)^i \hat{c}_{i,\uparrow} \hat{c}_{i,\downarrow}$ ,  $\hat{\eta}^+ = \hat{\eta}^{-\dagger}$ ,  $\hat{\eta}^z = \sum_i (\hat{n}_{i,\uparrow} + \hat{n}_{i,\downarrow} - 1)/2$

- $\hat{\eta}^+ |\text{vac}\rangle$  — doublon w/ quasimomentum  $\pi$  ( $\eta$  pair) — long-range coherence
- At half filling,  $\hat{\eta}^2 = \hat{\eta}^+ \hat{\eta}^-$  counts  $\eta$  pairs — eigenvals  $N_\eta(N_\eta + 1)$ ,  $N_\eta = 0, 1, \dots, l$
- Neighbouring sectors coupled by current operator  $\hat{\mathcal{J}}$  (imaginary tunnelling)

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Starting from ground state ( $N_\eta = 0$ ), apply  $\hat{\mathcal{J}}$  when it will increase  $\langle \hat{\eta}^2 \rangle$  the most

